

# Rigidity of Eigenvalues for $\beta$ Ensemble in Multi-cut Regime

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## Definition and some background of the $\beta$ ensemble

A  $\beta$  ensemble is a probability measure on  $\mathbb{R}^N$ :

$$\mu_N = \frac{1}{Z} \cdot e^{-\frac{N\beta}{2}[V(x_1)+\dots+V(x_N)]} \prod_{1 \leq i < j \leq N} |x_i - x_j|^\beta dx_1 \cdots dx_N$$

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- ▶  $Z$  is the normalization constant.

## Definition and some background of the $\beta$ ensemble

**Interest:** the behavior of  $x_1, \dots, x_N$  as  $N \rightarrow +\infty$ .

## Definition and some background of the $\beta$ ensemble

Recall the **Gaussian Orthogonal Ensemble (GOE)**:

$$M_N = \begin{pmatrix} \sqrt{2} a_{11} & a_{12} & \cdots & a_{1N} \\ a_{12} & \sqrt{2} a_{22} & \cdots & a_{2N} \\ \vdots & \vdots & \ddots & \vdots \\ a_{1N} & a_{2N} & \cdots & \sqrt{2} a_{NN} \end{pmatrix}$$

where  $a_{11}, a_{12}, a_{22}, a_{13}, \dots$  are i.i.d. standard normal random variables.

Since  $M_N$  is symmetric, all its eigenvalues are real.

# Definition and some background of the $\beta$ ensemble

## Theorem

If  $\beta = 1$  and  $V(x) = \frac{x^2}{2}$ , then  $\mu_N$  is the law of eigenvalues of  $\frac{1}{\sqrt{N}}M_N$ .

More precisely, if  $\lambda_1 \leq \dots \leq \lambda_N$  are eigenvalues of  $\frac{1}{\sqrt{N}}M_N$ , then

$$\mathbb{P}((\lambda_1, \dots, \lambda_N) \in \mathbf{A}) = N! \cdot \mu_N(\mathbf{A}) \quad \forall \mathbf{A} \subset \mathbb{R}^N \cap \{x_1 \leq \dots \leq x_N\}.$$



# Definition and some background of the $\beta$ ensemble



$$\mu_N \xrightarrow{\beta=1} \text{random symmetric matrix} \xrightarrow{V(x)=\frac{x^2}{2}} \text{GOE}$$

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$$\mu_N \xrightarrow{\beta=2} \text{random Hermitian matrix} \xrightarrow{V(x)=\frac{x^2}{2}} \text{GUE}$$

(GUE: **Gaussian Unitary Ensemble**)

# Definition and some background of the $\beta$ ensemble



$\mu_N \xrightarrow{\beta=1}$  random symmetric matrix  $\xrightarrow{V(x)=\frac{x^2}{2}}$  GOE

$\mu_N \xrightarrow{\beta=2}$  random Hermitian matrix  $\xrightarrow{V(x)=\frac{x^2}{2}}$  GUE

(GUE: Gaussian Unitary Ensemble)

$\mu_N \xrightarrow{\beta=4}$  random Hermitian quaternionic matrix  $\xrightarrow{V(x)=\frac{x^2}{2}}$  GSE

(GSE: Gaussian Symplectic Ensemble)

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$\mu_N \xrightarrow{V(x)=ax^2+bx+c}$  random tridiagonal matrix

(Dumitriu, Edelman (2002))

## Equilibrium measure of $\mu_N$

For  $(x_1, \dots, x_N) \in \mathbb{R}^N$ , define the empirical measure

$$L_N = \frac{1}{N} \sum_{i=1}^N \delta(t - x_i).$$

For example,  $\int_{\mathbb{R}} \sin t \, dL_N(t) = \frac{1}{N} \sum_{i=1}^N \sin x_i$ .

### Theorem

*Under some technical conditions,  $L_N$  converges to a limit  $\rho(t)dt$  in the following sense:*

$$\mathbb{E}^{\mu_N} \left[ \int_{\mathbb{R}} f(t) \, dL_N(t) \right] \xrightarrow{N \rightarrow +\infty} \int_{\mathbb{R}} f(t) \rho(t) \, dt$$

*for all bounded continuous function  $f(t)$ .*

*We call  $\rho(t)dt$  the **equilibrium measure**.*

## Equilibrium measure of $\mu_N$

Fact:

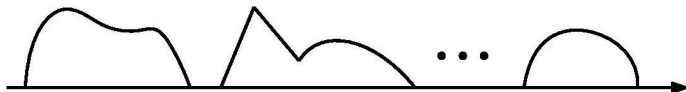
- ▶  $\rho(t)$  is continuous.

# Equilibrium measure of $\mu_N$

Fact:

- ▶  $\rho(t)$  is continuous.
- ▶  $\rho(t)$  has a compact support which is the union of finite many intervals:

$$\text{supp}(\rho) = [A_1, B_1] \cup \cdots \cup [A_q, B_q].$$

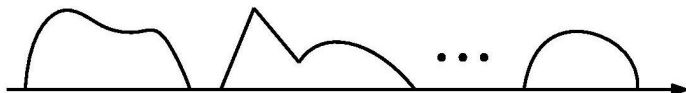


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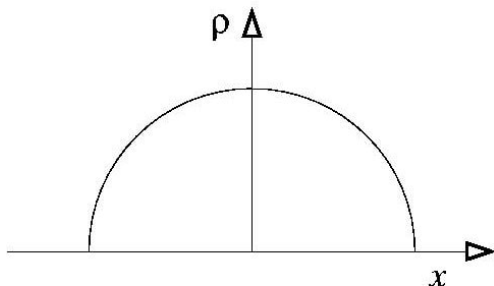
If  $q = 1$ , then we say that the  $\beta$  ensemble is in **one-cut regime**.  
If  $q > 1$ , then we say that the  $\beta$  ensemble is in **multi-cut regime**.  
Each  $[A_i, B_i]$  is called a **cut**.



## Equilibrium measure of $\mu_N$

For example, if  $\beta = 1$  and  $V(x) = \frac{x^2}{2}$ , then  $\rho(x)dx$  is the semicircle distribution:

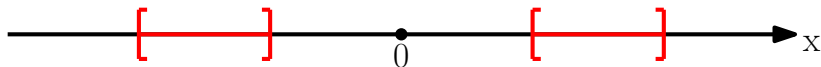
$$\rho(x) = \frac{1}{2\pi} \sqrt{4 - x^2} \mathbf{1}_{|x| \leq 2}$$



## Equilibrium measure of $\mu_N$

For example, if  $\beta > 0$  and  $V(x) = \frac{1}{2}x^4 - \frac{3\sqrt{2}\beta}{2}x^2$ , then  $\rho(x)dx$  is supported on:

$$\left[-\sqrt{2\sqrt{2}\beta}, -\sqrt{\sqrt{2}\beta}\right] \cup \left[\sqrt{\sqrt{2}\beta}, \sqrt{2\sqrt{2}\beta}\right]$$

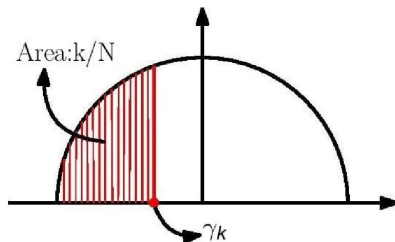


## Rigidity of particles for $\beta$ ensemble in one-cut regime

Under  $\mu_N$ ,  $(x_1, \dots, x_N)$  is a random element in  $\mathbb{R}^N$ . So each coordinate is a random element in  $\mathbb{R}$ . We call each  $x_i$  a **particle** or an **eigenvalue**.

**"classical location"** of the  $k$ -th largest particle:

$$\gamma_k = \gamma_k(N) = \inf_{t \in \mathbb{R}} \int_{-\infty}^t \rho(s) ds = \frac{k}{N} \quad k = 1, \dots, N$$



## Rigidity of particles for $\beta$ ensemble in one-cut regime

- ▶ Assume that the  $\beta$  ensemble is in the one-cut regime.
- ▶ For each  $x = (x_1, \dots, x_N)$ , use  $\bar{x}_k$  to denote its  $k$ -th largest particle.

Rigidity for particles in the bulk:

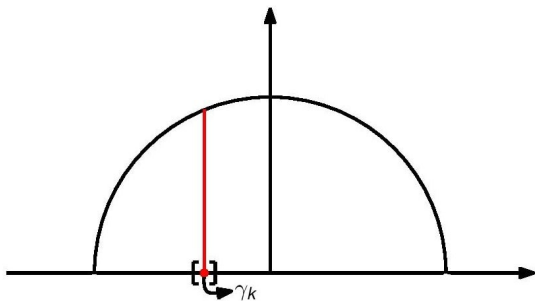
Theorem (Bourgade-Erdős-Yau (2011))

*Suppose  $a > 0$  and  $\epsilon > 0$ . Then under some technical conditions, there exists  $C > 0$  such that*

$$\mathbb{P}^{\mu_N} \left( |\bar{x}_k - \gamma_k| \leq N^{-1+\epsilon} \text{ for all } k \in [aN, (1-a)N] \right) \geq 1 - e^{-N^C}$$

*when  $N$  is large enough.*

## Rigidity of particles for $\beta$ ensemble in one-cut regime



$\bar{x}_k$ , i.e., the  $k$ -th largest particle, is basically located in an interval of size  $N^{-1}$  around  $\gamma_k$ .

# Rigidity of particles for $\beta$ ensemble in one-cut regime

## Rigidity for particles in the edge:

### Theorem (Bourgade-Erdős-Yau (2013))

Suppose  $\epsilon > 0$ . Then under some technical conditions, there exists  $C > 0$  such that

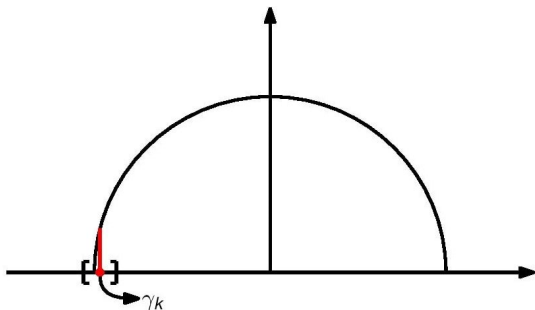
$$\mathbb{P}^{\mu_N} \left( |\bar{X}_k - \gamma_k| \leq N^{-\frac{2}{3} + \epsilon} \cdot \hat{k}^{-\frac{1}{3}} \text{ for all } k \in [1, N] \right) \geq 1 - e^{-N^C}$$

when  $N$  is large enough. Here  $\hat{k} := \min(k, N + 1 - k)$ .

### Remark

- ▶ When  $k \in [aN, (1 - a)N]$  for some  $a > 0$ , then  $N^{-\frac{2}{3} + \epsilon} \cdot \hat{k}^{-\frac{1}{3}} \sim N^{-1 + \epsilon}$ .
- ▶ When  $k$  is close to 1 or close to  $N$ , then  $N^{-\frac{2}{3} + \epsilon} \cdot \hat{k}^{-\frac{1}{3}} \sim N^{-\frac{2}{3} + \epsilon}$ .

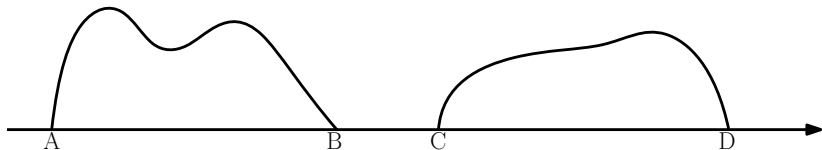
## Rigidity of particles for $\beta$ ensemble in one-cut regime



If  $k$  is close to 1 or close to  $N$ , i.e., near the edge, then the  $k$ -th largest particle  $\bar{x}_k$  is basically located in an interval of size  $N^{-2/3}$  around  $\gamma_k$ .

## Rigidity of particles for $\beta$ ensemble in multi-cut regime

The equilibrium measure  $\rho(x)dx$  is supported on more than one intervals:



For convenience, suppose  $\int_A^B \rho(x)dx = \int_C^D \rho(x)dx = \frac{1}{2}$ .



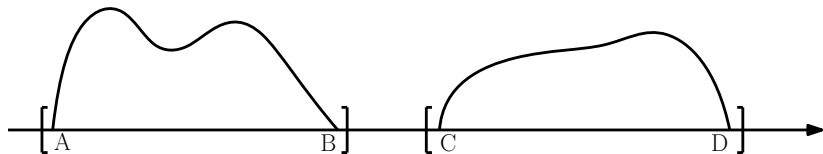
# Rigidity of particles for $\beta$ ensemble in multi-cut regime

## Theorem

Suppose  $\epsilon > 0$ . Under some technical conditions, there exists  $\delta > 0$  such that

$$\mathbb{P}^{\mu_N} \left( \text{all particles are in } [A - \epsilon, B + \epsilon] \cup [C - \epsilon, D + \epsilon] \right) \geq 1 - e^{-N^\delta}$$

for  $N$  large enough.



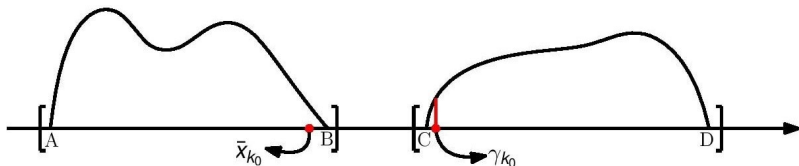
## Rigidity of particles for $\beta$ ensemble in multi-cut regime

Suppose  $k_0 = k_0(N) = \lfloor \frac{N}{2} \rfloor + 1$ . Then  $\gamma_{k_0}$  is the first classical location in  $[C, D]$ .

**Theorem (Bekerman (2015))**

*Suppose  $\epsilon > 0$ . Under some technical conditions we have that*

$$\overline{\lim} \mathbb{P}^{\mu_N} \left( \bar{x}_{k_0} \in [A - \epsilon, B + \epsilon] \right) := C_V > 0 \quad \text{as } N \rightarrow \infty.$$



Therefore, there is no rigidity near  $B$  and  $C$ .

# Rigidity of particles for $\beta$ ensemble in multi-cut regime

Rigidity for particles in the bulk:

Theorem (Li (2016))

*Suppose  $a > 0$  and  $\epsilon > 0$ . Then under some technical conditions, there exist  $C > 0$  such that*

$$\mathbb{P}^{\mu_N} \left( |\bar{x}_k - \gamma_k| \leq N^{-1+\epsilon} \text{ for all } k \in [aN, (\frac{1}{2} - a)N] \right) \geq 1 - e^{-N^C}$$

*when  $N$  is large enough. Same conclusion holds for  $k \in [(\frac{1}{2} + a)N, (1 - a)N]$*

## Idea of the proof of bulk rigidity for $\beta$ ensemble in multi-cut regime

- ▶ We cannot directly use the method of Bourgade, Erdős and Yau.
- ▶ Their first step is to show that for  $\epsilon > 0$ , there exist  $C > 0$  such that

$$\mathbb{P}^{\mu_N} \left( |\bar{x}_k - \gamma_k| \leq \epsilon \text{ for all } k \in [1, N] \right) \geq 1 - e^{-N^C}$$

when  $N$  is large enough.

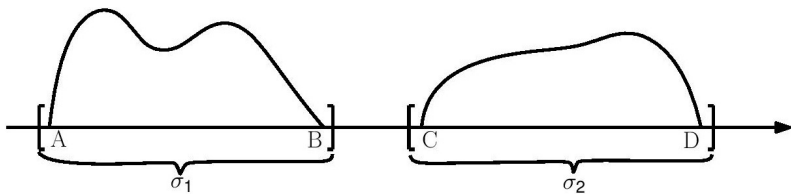
- ▶ But in multi-cut case particles may **jump**.

# Idea of the proof of bulk rigidity for $\beta$ ensemble in multi-cut regime

Idea: decompose the  $\beta$  ensemble in multi-cut regime as a product of  $\beta$  ensembles in one-cut regime.

Consider

- ▶  $\sigma_1$ : a neighbourhood of  $[A, B]$ ;
- ▶  $\sigma_2$ : a neighbourhood of  $[C, D]$



## Idea of the proof of bulk rigidity for $\beta$ ensemble in multi-cut regime

Recall that  $\mu_N = \frac{1}{Z} \cdot e^{-\frac{N\beta}{2}[V(x_1)+\dots+V(x_N)]} \prod_{i<j} |x_i - x_j|^\beta dx_1 \cdots dx_N$ .

## Idea of the proof of bulk rigidity for $\beta$ ensemble in multi-cut regime

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Define a probability measure on  $(\sigma_1 \cup \sigma_2)^N$ :

$$\tilde{\mu}_N = \frac{1}{\tilde{Z}} \cdot e^{-\frac{N\beta}{2} \sum_{i=1}^N V^r(x_i)} \cdot \prod_{\substack{i<j \\ x_i, x_j \in \sigma_1}} |x_i - x_j|^\beta \cdot \prod_{\substack{i<j \\ x_i, x_j \in \sigma_2}} |x_i - x_j|^\beta dx_1 \cdots dx_N$$

where  $\tilde{Z}$  is the normalization constant and

$$V^r(x) = \begin{cases} V_1(x) := V(x) - 2 \int_{[C,D]} \rho(y) \ln |x - y| dy & \text{if } x \in \sigma_1, \\ V_2(x) := V(x) - 2 \int_{[A,B]} \rho(y) \ln |x - y| dy & \text{if } x \in \sigma_2. \end{cases}$$

## Idea of the proof of bulk rigidity for $\beta$ ensemble in multi-cut regime

Fact:  $\tilde{\mu}_N$  is very close to  $\mu_N$ . If an event  $A_N$  is exponentially small under  $\tilde{\mu}_N$ , then it is also exponentially small under  $\mu_N$ :

$$\mathbb{P}^{\tilde{\mu}_N}(A_N) \geq 1 - e^{-N^C} \implies \mathbb{P}^{\mu_N}(A_N) \geq 1 - e^{-N^{C'}}$$

So we only have to show the rigidity with respect to  $\tilde{\mu}_N$ .



## Idea of the proof of bulk rigidity for $\beta$ ensemble in multi-cut regime

Under  $\tilde{\mu}_N$ , particles in different cuts don't have interaction.  
Basically,  $\tilde{\mu}_N = \eta_k \times \xi_s$  where

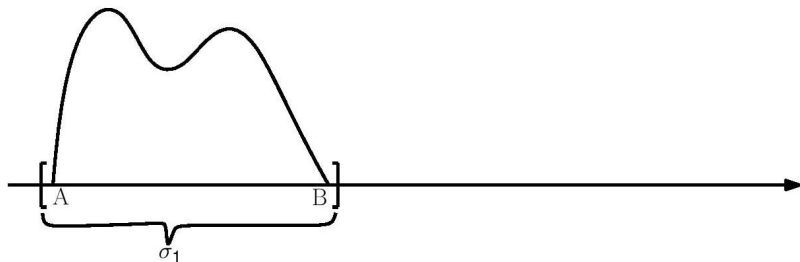
$$\eta_k = \frac{1}{Z_1} e^{-\frac{N\beta}{2} \sum_{i=1}^k V_1(x_i)} \prod_{1 \leq i < j \leq k} |x_i - x_j|^\beta \prod_{i=1}^k \mathbb{1}_{\sigma_1}(x_i) dx_1 \cdots dx_k$$

$$\xi_s = \frac{1}{Z_2} e^{-\frac{N\beta}{2} \sum_{i=1}^s V_2(x_i)} \prod_{1 \leq i < j \leq s} |x_i - x_j|^\beta \prod_{i=1}^s \mathbb{1}_{\sigma_2}(x_i) dx_1 \cdots dx_s$$

and  $k + s = N$

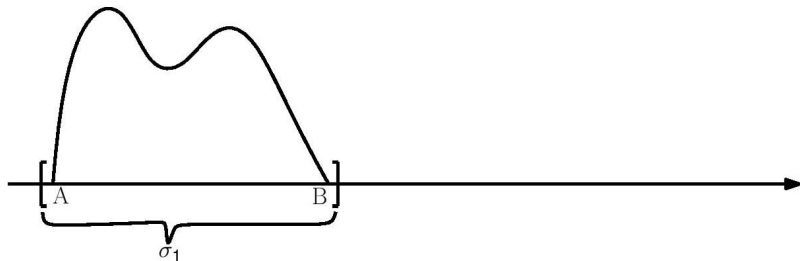
## Idea of the proof of bulk rigidity for $\beta$ ensemble in multi-cut regime

As  $k \rightarrow \infty$ , the equilibrium measure of  $\eta_k$  is  $2\rho(x)\mathbb{1}_{[A,B]}(x)dx$ :

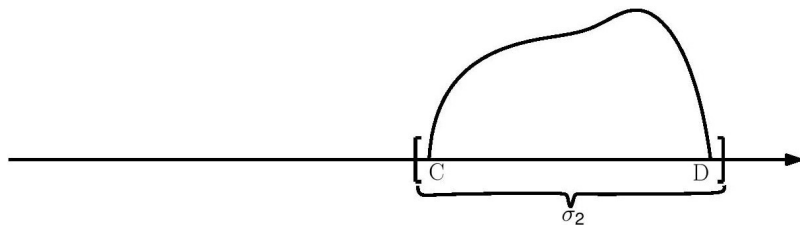


## Idea of the proof of bulk rigidity for $\beta$ ensemble in multi-cut regime

As  $k \rightarrow \infty$ , the equilibrium measure of  $\eta_k$  is  $2\rho(x)\mathbb{1}_{[A,B]}(x)dx$ :



As  $s \rightarrow \infty$ , the equilibrium measure of  $\xi_s$  is  $2\rho(x)\mathbb{1}_{[C,D]}(x)dx$ :



# Idea of the proof of bulk rigidity for $\beta$ ensemble in multi-cut regime

$\eta_k$  and  $\xi_s$  are in one-cut regime



rigidity for particles in the bulk with respect to  $\eta_k$  and  $\xi_s$



rigidity for particles in the bulk with respect to  $\tilde{\mu}_N = \eta_k \times \xi_s$



rigidity for particles in the bulk with respect to  $\mu_N$

# Conjecture: edge rigidity for $\beta$ ensemble in multi-cut regime

Conjecture:

$\eta_k$  and  $\xi_s$  are in one-cut regime



rigidity for particles near the **edges** with respect to  $\eta_k$  and  $\xi_s$



rigidity for particles near  $A$  and  $D$  with respect to  $\tilde{\mu}_N = \eta_k \times \xi_s$



rigidity for particles near  $A$  and  $D$  respect to  $\mu_N$

**Thank you!**