

The solution to the q -KdV equation*

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Abstract: Let KdV stand for the N th Gelfand-Dickey reduction of the KP hierarchy. The purpose of this paper is to show that any KdV solution leads effectively to a solution of the q -approximation of KdV. Two different q -KdV approximations were proposed, by E. Frenkel [7] and Khesin, Lyubashenko and Roger [12]. We show there is a dictionary between the solutions of q -KP and the 1-Toda lattice equations, obeying some special requirement; this is based on an algebra isomorphism between difference operators and D -operators, where $Df(x) = f(qx)$. Therefore every notion about the 1-Toda lattice can be transcribed into q -language. So, q -KdV is yet another Toda discretization of KdV.

Consider the q -difference operators D and D_q , defined by

$$Df(y) = f(qy) \quad \text{and} \quad D_q f(y) := \frac{f(qy) - f(y)}{(q-1)y},$$

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and the q -pseudo-differential operators

$$Q = D + u_0(x)D^0 + u_{-1}D^{-1} + \dots \quad \text{and} \quad Q_q = D_q + v_0(x)D_q^0 + v_{-1}(x)D_q^{-1} + \dots$$

The following q -versions of KP were proposed by E. Frenkel [7] and by Khesin, Lyubashenko and Roger [12], for $n = 1, 2, \dots$:

$$\frac{\partial Q}{\partial t_n} = [(Q^n)_+, Q] \quad (\text{Frenkel system}) \quad (0.1)$$

$$\frac{\partial Q_q}{\partial t_n} = [(Q_q^n)_+, Q_q], \quad (\text{KLR system}) \quad (0.2)$$

where $(\)_+$ and $(\)_-$ refer to the q -differential and strictly q -pseudo-differential part of $(\)$. The two systems are identical, after a (constant) upper-triangular linear transformation from the u_i 's to the v_i 's, as will become clear from the isomorphism between q -operators and difference operators, explained below. The purpose of this paper is to give a large class of solutions to both systems.

The δ -function $\delta(z) := \sum_{i \in \mathbf{Z}} z^i$; enjoys the property $f(\lambda, \mu)\delta(\lambda/\mu) = f(\lambda, \lambda)\delta(\lambda/\mu)$. Consider an appropriate space of functions $f(y)$ representable by "Fourier" series in the basis $\varphi_n(y) := \delta(q^{-n}x^{-1}y)$ for fixed $q \neq 1$,

$$f(y) = \sum_{-\infty}^{\infty} f_n \varphi_n(y);$$

the operators D , defined by $Df(y) = f(qy)$, and multiplication by a function $a(y)$ act on the basis elements, as follows:

$$D\varphi_n(y) = \varphi_{n-1}(y) \quad \text{and} \quad a(y)\varphi_n(y) = a(xq^n)\varphi_n(y).$$

Therefore, the Fourier transform,

$$f \longmapsto \hat{f} = (\dots, f_n, \dots)_{n \in \mathbf{Z}},$$

induces an algebra isomorphism, mapping D -operators onto a special class of Λ -operators in the shift $\Lambda := (\delta_{i,j-1})_{i,j \in \mathbf{Z}}$, as follows:

$$\sum_i a_i(y)D^i \longmapsto \sum_i \text{diag}(\dots, a_i(xq^n), \dots)_{n \in \mathbf{Z}} \Lambda^i; \quad (0.3)$$

conversely, any difference operator, depending on x , of the type (0.3) i.e., annihilated by $D - Ad_\Lambda$, where $(Ad_\Lambda)a = \Lambda a \Lambda^{-1}$, maps into a D -operator. This is the crucial basic isomorphism used throughout this paper.

To the shift Λ and to a fixed diagonal matrix $\lambda = \text{diag}(\lambda_{n+1})_{n \in \mathbf{Z}}$, we associate new operators

$$\tilde{\Lambda} = -\lambda\Lambda \quad \text{and} \quad \tilde{\tilde{\Lambda}} = \tilde{\Lambda} + \lambda = -\lambda(\Lambda - 1).$$

Observe that, under the isomorphism (0.3),

$$D \mapsto \Lambda, \quad \frac{1}{(q-1)x}D \mapsto \tilde{\Lambda} \quad \text{and} \quad D_q \mapsto \tilde{\tilde{\Lambda}},$$

upon setting $\lambda_n^{-1} = (1-q)xq^{n-1}$.

Defining the simple vertex operators

$$X(t, z) := e^{\sum_1^\infty t_i z^i} e^{-\sum_1^\infty \frac{z^{-i}}{i} \frac{\partial}{\partial t_i}}, \quad (0.4)$$

we now make a statement concerning the so-called one-Toda lattice; the latter describes deformations of a bi-infinite matrix L , which is lower-triangular, except for 1's just above the main diagonal. The *first formula* (0.6) below gives a solution to the *Frenkel system* (Theorem 0.1), upon replacing $\tilde{\Lambda}$ by Λ , which amounts to conjugating L by a constant diagonal matrix ε ; see (2.2). The *second formula* (0.6) gives, via the isomorphism (0.3), a solution to the *KLR system* (Theorem 0.2). Thus in the L -representation the two systems are related by a trivial diagonal conjugation. Note, Theorem 0.1 is given for arbitrary $\lambda = (\dots, \lambda_{-1}, \lambda_0, \lambda_1, \dots)$.

We shall need the well-known Hirota symbol for a polynomial p ,

$$p(\pm\tilde{\partial})f \circ g := p\left(\pm\frac{\partial}{\partial y_1}, \pm\frac{1}{2}\frac{\partial}{\partial y_2}, \dots\right) f(t+y)g(t-y)\Big|_{y=0}.$$

Note A_+ refers to the upper-triangular part of a matrix A , including the diagonal, and for $\alpha \in \mathbf{C}$, set $[\alpha] := (\alpha, \frac{1}{2}\alpha^2, \frac{1}{3}\alpha^3, \dots) \in \mathbf{C}^\infty$.

Theorem 0.1. *Given an integer $N \geq 2$, consider an arbitrary τ -function for the KP equation such that $\partial\tau/\partial t_{iN} = 0$ for $i = 1, 2, 3, \dots$ (N -KdV hierarchy).*

For a fixed $\lambda, \nu, c \in \mathbf{C}^\infty$, the infinite sequence of τ -functions ¹

$$\tau_n := X(t, \lambda_n) \dots X(t, \lambda_1) \tau(c + t), \quad \tau_0 = \tau(c + t), \quad \text{for } n \geq 0;$$

satisfies the 1-Toda bilinear identity for all $t, t' \in \mathbf{C}^\infty$ and all $n > m$:

$$\oint_{z=\infty} \tau_n(t - [z^{-1}]) \tau_{m+1}(t' + [z^{-1}]) e^{\sum_1^\infty (t_i - t'_i) z^i} z^{n-m-1} dz = 0.$$

The bi-infinite matrix (a full matrix below the main diagonal), where p_ℓ are the elementary Schur polynomials,

$$L = \sum_{\ell=0}^{\infty} \text{diag} \left(\frac{p_\ell(\tilde{\partial}) \tau_{n+2-\ell} \circ \tau_n}{\tau_{n+2-\ell} \tau_n} \right)_{n \in \mathbf{Z}} \tilde{\Lambda}^{1-\ell} \quad (0.5)$$

has the following properties:

(i) L^N satisfies the 1-Toda lattice

$$\frac{\partial L^N}{\partial t_n} = [(L^N)_+, L^N], \quad n = 1, 2, \dots,$$

(ii) L^N is upper triangular and admits the following expression in terms² of $\tilde{\Lambda}$ and $\tilde{\tilde{\Lambda}}$:

$$\begin{aligned} L^N &= \tilde{\Lambda}^N + \sum_1^N (\lambda_j + b_j) \tilde{\Lambda}^{N-1} + \left(\sum_0^{N-1} a_j + \sum_{1 \leq i \leq j \leq N-1} (\lambda_i + b_i)(\lambda_j + b_j) \right) \tilde{\Lambda}^{N-2} \\ &\quad + \dots + \lambda_1^N \tilde{\Lambda}^0 \\ &= \tilde{\tilde{\Lambda}}^N + \left(\sum_1^N b_j \right) \tilde{\tilde{\Lambda}}^{N-1} \\ &\quad + \left(\sum_0^{N-1} a_j - \sum_1^{N-1} (b_N - b_i) \lambda_i + \sum_{1 \leq i \leq j \leq N-1} b_i b_j \right) \tilde{\tilde{\Lambda}}^{N-2} + \dots \end{aligned} \quad (0.6)$$

¹ τ_n for $n < 0$ is defined later in (3.3).

²in the expressions below, the coefficients of the $\tilde{\tilde{\Lambda}}$'s are diagonal matrices, whose 0th component is given by the expression appearing below; e.g., $\sum_1^N b_j$ stands for $\text{diag}(\sum_1^N b_{j+n})_{n \in \mathbf{Z}}$ and λ_1^N stands for $\text{diag}(\lambda_{1+n}^N)_{n \in \mathbf{Z}}$.

with

$$b_k = \frac{\partial}{\partial t_1} \log \frac{\tau(c+t - \sum_1^k [\lambda_i^{-1}])}{\tau(c+t - \sum_1^{k-1} [\lambda_i^{-1}])}, \quad a_k = \left(\frac{\partial}{\partial t_1} \right)^2 \log \tau \left(c+t - \sum_1^k [\lambda_i^{-1}] \right), \quad (0.7)$$

for $k \geq 1$. These expressions for $k \leq 0$ will be given in (3.4) and (3.5).

In view of (0.7), the shift

$$\Lambda : b_k \longmapsto \Lambda b_k = b_{k+1} \quad \text{and} \quad a_k \longmapsto \Lambda a_k = a_{k+1}$$

corresponds to the following transformation,

$$\Lambda : c \mapsto c - [\lambda_1^{-1}] \quad \text{and} \quad \lambda_i \mapsto \lambda_{i+1}. \quad (0.8)$$

Therefore, in order that L^N satisfies the form of the right hand side of (0.3), we must make c and λ_i depend on x and q , such that the map Λ on a , b , λ corresponds to D , in addition to the fact that all λ_i must tend to ∞ simultaneously and c to $(x, 0, 0, \dots)$, when q goes to 1. So, $c(x)$ and $\lambda(x)$ must satisfy:

$$\begin{cases} Dc(x) = c(x) - [\lambda_1^{-1}(x)] \\ D\lambda_n(x) = \lambda_{n+1}(x) \\ \lim_{q \rightarrow 1} \lambda_i = \infty \\ \lim_{q \rightarrow 1} c(x) = \bar{x} := (x, 0, 0, \dots); \end{cases} \quad (0.9)$$

its only solution is given by:

$$\lambda_n^{-1} = (1-q)xq^{n-1} \quad \text{and} \quad c(x) = \left(\frac{(1-q)x}{1-q}, \frac{(1-q)^2 x^2}{2(1-q^2)}, \frac{(1-q)^3 x^3}{3(1-q^3)}, \dots \right), \quad (0.10)$$

and thus $D^n c(x) = c(x) - \sum_1^n [\lambda_i^{-1}]$. With this choice of λ_n ,

$$\frac{1}{(q-1)x} D \longmapsto \tilde{\Lambda} \quad \text{and} \quad D_q := \frac{D-1}{(q-1)x} \longmapsto \tilde{\tilde{\Lambda}}. \quad (0.11)$$

In analogy with (0.4), we define the simple q -vertex operators:

$$X_q(x, t, z) := e_q^{xz} X(t, z) \quad \text{and} \quad \tilde{X}_q(x, t, z) := \left(e_q^{xz} \right)^{-1} X(-t, z). \quad (0.12)$$

in terms of (0.4) and the q -exponential $e_q^x := e^{\sum_1^\infty \frac{(1-q)^k x^k}{k(1-q^k)}}$. Therefore under the isomorphism (0.3), Theorem 0.1 can be translated into q -language, to read:

Theorem 0.2. *Any KdV τ -function leads to a q -KdV τ -function $\tau(c(x)+t)$; the latter satisfies the bilinear relations below, for all $x \in \mathbf{R}$, $t, t' \in \mathbf{C}^\infty$, and all $n > m$, which tends to the standard KP-bilinear identity, when q goes to 1:*

$$\begin{aligned} \oint_{z=\infty} D^n (X_q(x, t, z) \tau(c(x) + t)) D^{m+1} (\tilde{X}_q(x, t', z) \tau(c(x) + t')) dz &= 0 \\ \longrightarrow \oint_{z=\infty} X(t, z) \tau(\bar{x} + t) X(t', z) \tau(\bar{x} + t') dz &= 0 \end{aligned} \quad (0.13)$$

Moreover, the q -differential operator Q_q^N has the form below and tends to the differential operator \mathcal{L}^N of the KdV hierarchy, when q goes to 1:

$$\begin{aligned} Q_q^N &= D_q^N + \frac{\partial}{\partial t_1} \log \frac{\tau(D^N c + t)}{\tau(c + t)} D_q^{N-1} \\ &\quad + \left(\sum_{i=0}^{N-1} \frac{\partial^2}{\partial t_1^2} \log \tau(D^i c + t) \right) \\ &\quad - \sum_{i=0}^{N-2} \lambda_{i+1} \left(\frac{\partial}{\partial t_1} \log \frac{\tau(D^N c + t)}{\tau(D^{N-1} c + t)} - \frac{\partial}{\partial t_1} \log \frac{\tau(D^{i+1} c + t)}{\tau(D^i c + t)} \right) \\ &\quad + \sum_{0 \leq i \leq j \leq N-2} \frac{\partial}{\partial t_1} \log \frac{\tau(D^{i+1} c + t)}{\tau(D^i c + t)} \frac{\partial}{\partial t_1} \log \frac{\tau(D^{j+1} c + t)}{\tau(D^j c + t)} D_q^{N-2} + \dots \\ &\longrightarrow \left(\frac{\partial}{\partial x} \right)^N + N \frac{\partial^2}{\partial t_1^2} \log \tau(\bar{x} + t) \left(\frac{\partial}{\partial x} \right)^{N-2} + \dots \end{aligned} \quad (0.14)$$

M.A. and PvM thank Edward Frenkel for kindly discussing this problem during spring 1996. For a systematic study of discrete systems, see Kupershmidt [13] and Gieseker [8]. It is an old observation (see [11]) that the Toda lattice discretizes the KdV equation, and this in many different ways. Therefore it is not surprising that q -KdV is yet another Toda discretization of KdV. In an elegant recent preprint, Iliev [10] has obtained q -bilinear identities and q -tau functions, as well, purely within the KP theory.

1 The KP τ -functions and Grassmannians

KP τ -functions satisfy the differential Fay identity for all $y, z \in \mathbf{C}$, in terms of the Wronskian $\{f, g\} := f'g - fg'$, as shown in [1, 16]:

$$\begin{aligned} & \{\tau(t - [y^{-1}]), \tau(t - [z^{-1}])\} \\ & + (y - z)(\tau(t - [y^{-1}])\tau(t - [z^{-1}]) - \tau(t)\tau(t - [y^{-1}] - [z^{-1}])) = 0. \end{aligned} \quad (1.1)$$

In fact this identity characterizes the τ -function, as shown in [15]. We shall need the following, shown in [1]:

Proposition 1.1. *Consider τ -functions τ_1 and τ_2 , the corresponding wave functions*

$$\Psi_i = e^{\sum_{i \geq 1} t_i z^i} \frac{\tau_i(t - [z^{-1}])}{\tau_i(t)} = e^{\sum_{i \geq 1} t_i z^i} (1 + O(z^{-1})) \quad (1.2)$$

and the associated infinite-dimensional planes, as points in the Grassmannian Gr ,

$$\tilde{W}_i = \text{span} \left\{ \left(\frac{\partial}{\partial t_1} \right)^k \Psi_i(t, z), \text{ for } k = 0, 1, 2, \dots \right\};$$

then the following statements are equivalent

- (i) $z\tilde{W}_2 \subset \tilde{W}_1$;
- (ii) $z\Psi_2(t, z) = \frac{\partial}{\partial t_1} \Psi_1(t, z) - \alpha\Psi_1(t, z)$, for some function $\alpha = \alpha(t)$;
- (iii)

$$\{\tau_1(t - [z^{-1}]), \tau_2(t)\} + z(\tau_1(t - [z^{-1}])\tau_2(t) - \tau_2(t - [z^{-1}])\tau_1(t)) = 0 \quad (1.3)$$

When (i), (ii) or (iii) holds, $\alpha(t)$ is given by

$$\alpha(t) = \frac{\partial}{\partial t_1} \log \frac{\tau_2}{\tau_1}. \quad (1.4)$$

Proof: To prove that (i) \Rightarrow (ii), the inclusion $z\tilde{W}_2 \subset \tilde{W}_1$ implies $z\tilde{W}_2^t \subset \tilde{W}_1^t$, where $\tilde{W}_1^t = \tilde{W}_1 e^{-\sum_{i \geq 1} t_i z^i}$; it follows that

$$z\psi_2(t, z) = z(1 + O(z^{-1})) \in W_1^t$$

must be a linear combination, involving the wave functions $\Psi_i = e^{\sum_1^\infty t_i z^i} \psi_i$:

$$z\psi_2 = \left(\frac{\partial}{\partial x} + z \right) \psi_1 - \alpha(t)\psi_1, \text{ and thus } z\Psi_2 = \frac{\partial}{\partial t_1} \Psi_1 - \alpha(t)\Psi_1. \quad (1.5)$$

The expression (1.4) for $\alpha(t)$ follows from equating the z^0 -coefficient in (1.5), upon using the τ -function representation (1.2). To show that (ii) \Rightarrow (i), note that

$$z\Psi_2 = \frac{\partial}{\partial t_1} \Psi_1 - \alpha\Psi_1 \in W_1^0$$

and taking z -derivatives, we have

$$z \left(\frac{\partial}{\partial t_1} \right)^j \Psi_2 = \left(\frac{\partial}{\partial t_1} \right)^{j+1} \Psi_1 + \beta_1 \left(\frac{\partial}{\partial t_1} \right)^j \Psi_1 + \cdots + \beta_{j+1} \Psi_1,$$

for some $\beta_1, \dots, \beta_{j+1}$ depending on t only; this implies the inclusion (i). The equivalence (ii) \iff (iii) follows from a straight forward computation using the τ -function representation (1.2) of (ii) and the expression for $\alpha(t)$. \blacksquare

2 The one-Toda lattice

For details on this sketchy exposition, see [3]. The one-Toda lattice equations

$$\frac{\partial L}{\partial t_n} = [(L^n)_+, L], \quad (2.1)$$

are deformations of an infinite matrix

$$L = \sum_{-\infty < i \leq 0} a_i \tilde{\Lambda}^i + \tilde{\Lambda}, \text{ with } \tilde{\Lambda} := \lambda \Lambda = \varepsilon \Lambda \varepsilon^{-1}, \quad (2.2)$$

for diagonal matrices λ and ε , with non-zero entries, and diagonal matrices a_i , depending on $t = (t_1, t_2, \dots)$. Note the conjugation by the constant diagonal matrix ε is harmless, but it is necessary to capture the KLR-system. One introduces wave and adjoint wave vectors $\Psi(t, z)$ and $\Psi^*(t, z)$, satisfying

$$L\Psi = z\Psi \quad \text{and} \quad L^\top \Psi^* = z\Psi^*$$

and

$$\frac{\partial \Psi}{\partial t_n} = (L^n)_+ \Psi \quad \frac{\partial \Psi^*}{\partial t_n} = -((L^n)_+)^{\top} \Psi^*. \quad (2.3)$$

The wave vectors Ψ and Ψ^* can be expressed in terms of one sequence of τ -functions $\tau(n, t) := \tau_n(t_1, t_2, \dots)$, $n \in \mathbf{Z}$, to wit:

$$\begin{aligned} \Psi(t, z) &= \left(e^{\sum_1^{\infty} t_i z^i} \psi(t, z) \right)_{n \in \mathbf{Z}} = \left(\frac{\tau_n(t - [z^{-1}])}{\tau_n(t)} e^{\sum_1^{\infty} t_i z^i} \varepsilon_n z^n \right)_{n \in \mathbf{Z}}, \\ \Psi^*(t, z) &= \left(e^{-\sum_1^{\infty} t_i z^i} \psi^*(t, z) \right)_{n \in \mathbf{Z}} = \left(\frac{\tau_{n+1}(t + [z^{-1}])}{\tau_{n+1}(t)} e^{-\sum_1^{\infty} t_i z^i} \varepsilon_n^{-1} z^{-n} \right)_{n \in \mathbf{Z}} \end{aligned} \quad (2.4)$$

It follows that, in terms of $\chi(z) := (z^n)_{n \in \mathbf{Z}}$ and the notation $a_{\Lambda} := \text{diag}(a_{k+1})_{k \in \mathbf{Z}}$:

$$\Psi = e^{\sum_1^{\infty} t_i z^i} S \varepsilon \chi(z), \quad \text{with} \quad S = \sum_0^{\infty} \frac{p_n(-\tilde{\partial}) \tau(t)}{\tau(t)} \tilde{\Lambda}^{-n},$$

$$\Psi^* = e^{-\sum_1^{\infty} t_i z^i} (S^{\top})^{-1} \varepsilon^{-1} \chi(z^{-1}), \quad \text{with} \quad S^{-1} = \sum_0^{\infty} \tilde{\Lambda}^{-n} \left(\frac{p_n(\tilde{\partial}) \tau(t)}{\tau(t)} \right)_{\Lambda}.$$

Moreover, as will follow from Proposition 2.1 below, Ψ and Ψ^* satisfy the bilinear identities:

$$\oint_{z=\infty} \Psi_n(t, z) \Psi_m^*(t', z) \frac{dz}{2\pi i z} = 0, \quad \text{for all } n > m.$$

From the representation of S and S^{-1} above, it follows that

$$\begin{aligned} L^k &= S \tilde{\Lambda}^k S^{-1} \\ &= \sum_{\ell=0}^{\infty} \text{diag} \left(\frac{p_{\ell}(\tilde{\partial}) \tau_{n+k-\ell+1} \circ \tau_n}{\tau_{n+k-\ell+1} \tau_n} \right)_{n \in \mathbf{Z}} \tilde{\Lambda}^{k-\ell} \\ &= \tilde{\Lambda}^k + \text{diag} \left(\frac{\partial}{\partial t_1} \log \frac{\tau_{n+k}}{\tau_n} \right)_{n \in \mathbf{Z}} \tilde{\Lambda}^{k-1} + \dots \\ &\quad + \text{diag} \left(\frac{\partial}{\partial t_k} \log \frac{\tau_{n+1}}{\tau_n} \right)_{n \in \mathbf{Z}} \tilde{\Lambda}^0 + \text{diag} \left(\frac{\partial^2}{\partial t_1 \partial t_k} \log \tau_n \right)_{n \in \mathbf{Z}} \tilde{\Lambda}^{-1} + \dots \end{aligned} \quad (2.5)$$

For instance, the Λ^0 -term in the last expression follows from setting $m = n - 2$, $t \mapsto t + [\alpha]$, $t' \mapsto t - [\alpha]$ in the bilinear identity above, yielding

$$\begin{aligned} 0 &= \frac{\tau_n(t + [\alpha])\tau_{n-1}(t - [\alpha])}{\tau_n(t)\tau_{n-1}(t)} \oint_{z=\infty} \Psi_n(t + [\alpha], z)\Psi_{n-2}^*(t - [\alpha], z)\frac{dz}{2\pi iz} \\ &= \frac{1}{\tau_n\tau_{n-1}} \sum_{j \geq 0} \alpha^j \left(\frac{\partial}{\partial t_{j+2}} - p_{j+2}(\tilde{\partial}) \right) \tau_n \circ \tau_{n-1}, \end{aligned}$$

and thus,

$$\frac{\partial}{\partial t_k} \log \frac{\tau_{n+1}}{\tau_n} = \frac{p_k(\tilde{\partial})\tau_{n+1} \circ \tau_n}{\tau_{n+1}\tau_n} = (L^k)_{nn}.$$

With each component of the wave vector Ψ , we associate a sequence of infinite-dimensional planes in the Grassmannian Gr

$$\begin{aligned} W_n &= \text{span}_{\mathbf{C}} \left\{ \left(\frac{\partial}{\partial t_1} \right)^k \Psi_n(t, z), \quad k = 0, 1, 2, \dots \right\} \\ &= e^{\sum_1^\infty t_i z^i} \text{span}_{\mathbf{C}} \left\{ \left(\frac{\partial}{\partial t_1} + z \right)^k \psi_n(t, z), \quad k = 0, 1, 2, \dots \right\} \quad (2.6) \end{aligned}$$

and planes

$$W_n^* = \text{span}_{\mathbf{C}} \left\{ \left(\frac{\partial}{\partial t_1} \right)^k \Psi_{n-1}^*(t, z), \quad k = 0, 1, 2, \dots \right\},$$

which are orthogonal to W_n by the residue pairing

$$\oint_{z=\infty} f(z)g(z)\frac{dz}{2\pi iz}. \quad (2.7)$$

Note that the plane $z^{-n}W_n$ has so-called virtual genus zero, in the terminology of [14]; in particular, this plane contains an element of order $1 + O(z^{-1})$. The following statement is contained in [3]:

Proposition 2.1. *The following five statements are equivalent*
(i) *The 1-Toda lattice equations (2.1)*

(ii) Ψ and Ψ^* , with the proper asymptotic behaviour, given by (2.4), satisfy the bilinear identities for all $t, t' \in \mathbf{C}^\infty$

$$\oint_{z=\infty} \Psi_n(t, z) \Psi_m^*(t', z) \frac{dz}{2\pi iz} = 0, \quad \text{for all } n > m; \quad (2.8)$$

(iii) the τ -vector satisfies the following bilinear identities for all $n > m$ and $t, t' \in \mathbf{C}^\infty$:

$$\oint_{z=\infty} \tau_n(t - [z^{-1}]) \tau_{m+1}(t' + [z^{-1}]) e^{\sum_1^\infty (t_i - t'_i) z^i} z^{n-m-1} dz = 0; \quad (2.9)$$

(iv) The components τ_n of a τ -vector correspond to a flag of planes in Gr ,

$$\supset W_{n-1} \supset W_n \supset W_{n+1} \supset \dots \quad (2.10)$$

(v) A sequence of KP- τ -functions τ_n satisfying the equations

$$\{\tau_n(t - [z^{-1}]), \tau_{n+1}(t)\} + z(\tau_n(t - [z^{-1}])\tau_{n+1}(t) - \tau_{n+1}(t - [z^{-1}])\tau_n(t)) = 0 \quad (2.11)$$

Proof: The proof that (i) is equivalent to (ii) follows from the methods in [4, 16]. That (ii) is equivalent to (iii) follows from the representation (2.4) of wave functions in terms of τ -functions. Finally, we sketch the proof that (ii) is equivalent to (iv). The inclusion in (iv) implies that W_n , given by (2.6), is also given by

$$W_n = \text{span}_{\mathbf{C}} \{\Psi_n(t, z), \Psi_{n+1}(t, z), \dots\};$$

Since each τ_n is a τ -function, we have that

$$\oint_{z=\infty} \Psi_n(t, z) \Psi_{n-1}^*(t', z) \frac{dz}{2\pi iz} = 0,$$

implying that, for each $n \in \mathbf{Z}$, $\Psi_{n-1}^*(t, z) \in W_n^*$. Moreover the inclusions $\dots \supset W_n \supset W_{n+1} \supset \dots$ imply, by orthogonality, the inclusions $\dots \subset W_n^* \subset W_{n+1}^* \subset \dots$, and thus

$$W_n^* = \{\Psi_{n-1}^*(t, z), \Psi_{n-2}^*(t, z), \dots\}.$$

Since

$$W_n \subset W_m = (W_m^*)^*, \quad \text{all } n \geq m,$$

we have the orthogonality $W_n \perp W_m^*$ by the residue pairing (2.7) for all $n \geq m$, i.e.,

$$\oint_{z=\infty} \Psi_n(t, z) \Psi_{m-1}^*(t', z) \frac{dz}{2\pi iz} = 0, \quad \text{all } n \geq m.$$

Note (ii) implies $W_m^* \subset W_n^*$, $n > m$, hence $W_n \subset W_m$, $n > m$, yielding (iv). That (iv) \iff (v) follows from proposition 1.1, by setting $\tau_1 := \tau_n$ and $\tau_2 = \tau_{n+1}$. Then (v) is equivalent to the inclusion property

$$z(z^{-n-1}W_{n+1}) \subset (z^{-n}W_n), \quad \text{i.e. } W_{n+1} \subset W_n,$$

thus ending the proof of proposition 2.1. ■

3 Proof of Theorems 0.1 and 0.2

At first, we exhibit particular solutions to equation (2.11), explained in [1].

Lemma 3.1. *Particular solutions to equation*

$$\{\tau_1(t - [z^{-1}]), \tau_2(t)\} + z(\tau_1(t - [z^{-1}])\tau_2(t) - \tau_2(t - [z^{-1}])\tau_1(t)) = 0$$

are given, for arbitrary $\lambda \in \mathbf{C}^*$, by pairs (τ_1, τ_2) , defined by:

$$\tau_2(t) = X(t, \lambda)\tau_1(t) = e^{\sum t_i \lambda^i} \tau_1(t - [\lambda^{-1}]), \quad (3.1)$$

or

$$\tau_1(t) = X(-t, \lambda)\tau_2(t) = e^{-\sum t_i \lambda^i} \tau_2(t + [\lambda^{-1}]). \quad (3.2)$$

Proof: Using

$$e^{-\sum_1^\infty \frac{1}{i} (\frac{\lambda}{z})^i} = 1 - \frac{\lambda}{z},$$

it suffices to check that $\tau_2(t)$ satisfies the above equation (2.11)

$$\begin{aligned} & e^{-\sum t_i \lambda^i} \left(\{\tau_1(t - [z^{-1}]), \tau_2(t)\} + z(\tau_1(t - [z^{-1}])\tau_2(t) - \tau_2(t - [z^{-1}])\tau_1(t)) \right) \\ & = e^{-\sum t_i \lambda^i} \{ \tau_1(t - [z^{-1}]), e^{\sum t_i \lambda^i} \tau_1(t - [\lambda^{-1}]) \} \end{aligned}$$

$$\begin{aligned}
& + z(\tau_1(t - [z^{-1}])\tau_1(t - [\lambda^{-1}]) - (1 - \frac{\lambda}{z})\tau_1(t)\tau_1(t - [z^{-1}] - \lambda^{-1})) \\
= & \{\tau_1(t - [z^{-1}]), \tau_1(t - [\lambda^{-1}])\} \\
& + (z - \lambda)(\tau_1(t - [z^{-1}])\tau_1(t - [\lambda^{-1}]) - \tau_1(t)\tau_1(t - [z^{-1}] - [\lambda^{-1}])) \\
= & 0,
\end{aligned}$$

using the differential Fay identity (1.1) for the τ -function τ_1 ; a similar proof works for the second solution, given by (3.2). \blacksquare

Proof of Theorems 0.1 and 0.2: From an arbitrary N -KdV τ -function, construct, for $\lambda, c, \nu \in \mathbf{C}^\infty$, the following sequence of τ -functions, for $n \geq 0$, as announced in Theorem 0.1:

$$\begin{aligned}
\tau_0(t) &= \tau(c + t) \\
\tau_n &= X(t, \lambda_n) \dots X(t, \lambda_1) \tau(c + t) \\
&= \frac{\Delta(\lambda_1, \dots, \lambda_n)}{\prod_1^n \lambda_i^{i-1}} \prod_{k=1}^n e^{\sum_{i=1}^\infty t_i \lambda_k^i} \tau(c + t - \sum_1^n [\lambda_i^{-1}]), \\
\tau_{-n} &= X(-t, \lambda_{-n+1}) \dots X(-t, \lambda_0) \tau(c + t) \\
&= \frac{\Delta(\lambda_0, \dots, \lambda_{-n+1})}{\prod_1^n \lambda_{-i+1}^{i-1}} \prod_{k=1}^n e^{-\sum_{i=1}^\infty t_i \lambda_{-k+1}^i} \tau(c + t + \sum_1^n [\lambda_{-i+1}^{-1}])
\end{aligned} \tag{3.3}$$

and so, each τ_n is defined inductively by

$$\tau_{n+1} = X(t, \lambda_{n+1}) \tau_n;$$

thus by Lemma 3.1, the functions τ_{n+1} and τ_n are a solution of equation (v) of proposition 2.1. Therefore, by the same proposition 2.1, the τ_n 's form a τ -vector of the 1-Toda lattice. By removing the harmless exponential factor $\prod_{k=1}^n \exp(\sum_1^\infty t_{iN} \lambda_k^{iN})$, each τ_n has the property that $\partial \tau_n / \partial t_{iN} = 0$ for $i = 1, 2, \dots$; therefore

$$z^N W_n \subset W_n.$$

In particular, the representation

$$W_n = \text{span}\{\Psi_n(t, z), \Psi_{n+1}(t, z), \dots\},$$

which follows from the inclusion $\dots \supset W_n \supset W_{n+1} \supset \dots$, implies that, since $L\Psi = z\Psi$,

$$z^N \Psi_k = \sum_{j \geq k} a_j \Psi_j = (L^N \Psi)_k,$$

and thus L^N is upper-triangular.

Therefore, we conclude that the matrix L , defined by (2.5), from the sequence of τ -functions (3.3),

$$\begin{aligned} L &= \tilde{\Lambda} + \left(\frac{\partial}{\partial t_1} \log \frac{\tau_{n+1}}{\tau_n} \right)_{n \in \mathbf{Z}} + \left(\left(\frac{\partial}{\partial t_1} \right)^2 \log \tau_n \right)_{n \in \mathbf{Z}} \tilde{\Lambda}^{-1} + \dots \\ &= \tilde{\Lambda} + (\lambda_{n+1} + b_{n+1})_{n \in \mathbf{Z}} \tilde{\Lambda}^0 + (a_n)_{n \in \mathbf{Z}} \tilde{\Lambda}^{-1} + \dots, \end{aligned}$$

satisfies the 1-Toda lattice equations, where

$$\begin{aligned} b_{n+1} &= \frac{\partial}{\partial t_1} \log \frac{\tau(c+t - \sum_1^{n+1} [\lambda_i^{-1}])}{\tau(c+t - \sum_1^n [\lambda_i^{-1}])} \quad \text{for } n \geq 1 \\ &= \frac{\partial}{\partial t_1} \log \frac{\tau(c+t - [\lambda_1^{-1}])}{\tau(c+t)}, \quad \text{for } n = 0, \\ &= \frac{\partial}{\partial t_1} \log \frac{\tau(c+t + \sum_0^{n+2} [\lambda_i^{-1}](1 - \delta_{-1,n}))}{\tau(c+t + \sum_0^{n+1} [\lambda_i^{-1}])}, \quad \text{for } n \leq -1, \end{aligned} \quad (3.4)$$

and

$$\begin{aligned} a_n &= \frac{\partial^2}{\partial t_1^2} \log \tau(c+t - \sum_1^n [\lambda_i^{-1}]) \quad \text{for } n \geq 1 \\ &= \frac{\partial^2}{\partial t_1^2} \log \tau(c+t) \quad \text{for } n = 0 \\ &= \frac{\partial^2}{\partial t_1^2} \log \tau(c+t + \sum_0^{n+1} [\lambda_i^{-1}]) \quad \text{for } n \leq -1, \end{aligned} \quad (3.5)$$

confirming (0.7). Using the fact that, in view of (2.5), the diagonal terms of L^N are given by

$$\frac{\partial}{\partial t_N} \log \frac{\tau_{n+1}}{\tau_n} = \lambda_{n+1}^N,$$

and the fact that, in the notation of footnote 2,

$$\tilde{\Lambda}^n = (\tilde{\Lambda} + \lambda)^n = \tilde{\Lambda}^n + \left(\sum_1^n \lambda_i \right) \tilde{\Lambda}^{n-1} + \left(\sum_{1 \leq i \leq j \leq n-1} \lambda_i \lambda_j \right) \tilde{\Lambda}^{n-2} + \dots,$$

one finds that the upper-triangular matrix L^N has the following expression:

$$\begin{aligned}
L^N &= \tilde{\Lambda}^N + \sum_1^N (\lambda_j + b_j) \tilde{\Lambda}^{j-1} + \left(\sum_0^{N-1} a_j + \sum_{1 \leq i \leq j \leq N-1} (\lambda_i + b_i)(\lambda_j + b_j) \right) \tilde{\Lambda}^{N-2} \\
&\quad + \dots + \lambda_1^N \tilde{\Lambda}^0 \\
&= \tilde{\Lambda}^N + \left(\sum_1^N b_j \right) \tilde{\Lambda}^{N-1} + \left(\sum_0^{N-1} a_j - \sum_1^{N-1} (b_N - b_i) \lambda_i + \sum_{1 \leq i \leq j \leq N-1} b_i b_j \right) \tilde{\Lambda}^{N-2} + \dots
\end{aligned} \tag{3.6}$$

in terms of b_k and a_k defined in (0.7), thus proving Theorem 0.1. \blacksquare

To prove Theorem 0.2, note at first:

$$\begin{aligned}
\frac{z^{n-m-1}}{\prod_{k=m+2}^n (-\lambda_k)} \prod_{k=m+2}^n e^{-\sum_{i=1}^{\infty} \frac{1}{i} \left(\frac{\lambda_k}{z} \right)^i} &= \frac{z^{n-m-1}}{\prod_{k=m+2}^n (-\lambda_k)} \prod_{k=m+2}^n \left(1 - \frac{\lambda_k}{z} \right) \\
&= \prod_{k=m+2}^n \left(1 - \frac{z}{\lambda_k} \right) \\
&= \prod_{k=m+2}^n e^{-\sum_{i=1}^{\infty} \frac{1}{i} \left(\frac{z}{\lambda_k} \right)^i} \\
&= \frac{e^{xzq^n}}{e_q^{xzq^{m+1}}} = D^n e_q^{xz} D^{m+1} \left(e_q^{xz} \right)^{-1}.
\end{aligned}$$

The function τ_n , defined in Theorem 0.1, satisfies the bilinear identity of Theorem 0.1; therefore, using (3.3) and the above in the computation of $\tau_n(t - [z^{-1}])$, the following relations hold, up to a multiplicative factor depending on λ and ν :

$$\begin{aligned}
&\alpha(\lambda, \nu) \oint_{z=\infty} \tau_n(t - [z^{-1}]) \tau_{m+1}(t' + [z^{-1}]) e^{\sum_1^{\infty} (t_i - t'_i) z^i} z^{n-m} \frac{dz}{z} \\
&= \oint_{z=\infty} \tau(c(x) + t - [z^{-1}] - \sum_1^n [\lambda_i^{-1}]) \tau(c(x) + t' + [z^{-1}] + \sum_1^{m+1} [\lambda_i^{-1}]) \\
&\quad \prod_{k=m+2}^n \left(1 - \frac{z}{\lambda_k} \right) e^{\sum_1^{\infty} (t_i - t'_i) z^i} dz \\
&= \oint_{z=\infty} D^n (X_q(x, t, z) \tau(c(x) + t)) D^{m+1} (\tilde{X}_q(x, t', z) \tau(c(x) + t')) dz = 0.
\end{aligned}$$

When $q \rightarrow 1$, the second expression above tends to the standard KP-bilinear equation, upon using (0.10). Moreover, one checks by induction, using the expression (2.5) for L and (3.3), that $(L^N)_+$ for $N = 1, 2, 3, \dots$ has the q -form (0.3). Also, note that a_k and b_k can be expressed in terms of the D -operator, using (0.7); to wit:

$$b_k = \frac{\partial}{\partial t_1} \log \frac{\tau(D^k c + t)}{\tau(D^{k-1} c + t)}, \quad a_k = \left(\frac{\partial}{\partial t_1} \right)^2 \log \tau(D^k c + t).$$

So, the expression for Q_q^N in Theorem 0.2 follows at once from (3.6). The fact that

$$-\lambda_1 \frac{\partial}{\partial t_1} \log \frac{\tau(D^{j+1} c + t)}{\tau(D^j c + t)} \longrightarrow \frac{\partial^2}{\partial x^2} \log \tau(\bar{x} + t)$$

implies that all terms in (0.14) vanish in the limit $q \rightarrow 1$, except for the term $\sum_{i=0}^{N-1} \frac{\partial^2}{\partial t_1^2} \log \tau(D^i c + t)$; so we have that

$$\lim_{q \rightarrow 1} Q_q^N = \left(\frac{\partial}{\partial x} \right)^N + N \frac{\partial^2}{\partial x^2} \log \tau(\bar{x} + t) \left(\frac{\partial}{\partial x} \right)^{N-2} + \dots,$$

thus ending the proof of theorem 0.2. ■

4 Examples and vertex operators

The isomorphism (0.3) enables one to translate every 1-Toda statement, having the form (0.3) into a D or D_q statement. Also every τ -function of the KdV hierarchy leads automatically to a solution of q -KdV. For instance, by replacing $t \mapsto c(x) + t$ in the Schur polynomials, one finds q -Schur polynomials. The latter were obtained by Haine and Iliev [9] by using the q -Darboux transforms; the latter had been studied by Horozov and coworkers in [5, 6].

The n -soliton solution to the KdV (for $N = 2$) (for this formulation, see [4]),

$$\tau(t) = \det \left(\delta_{i,j} - \frac{a_j}{y_i + y_j} e^{-\sum_{k:\text{odd}} t_k (y_i^k + y_j^k)} \right)_{1 \leq i, j \leq n},$$

leads to a q -soliton by the shift $t \mapsto c(x) + t$, with $c(x)$ as in (0.8), namely

$$\tau(x, t) = \det \left(\delta_{ij} - a_j \frac{\left(\frac{e^{xy_i} e^{xy_i}}{q} \right)^{-1}}{y_i + y_j} e^{-\sum_{k=1}^{\infty} t_k (y_i^k + y_j^k)} \right)_{1 \leq i, j \leq n}.$$

Moreover the vertex operator for the 1-Toda lattice is a reduction of the 2-Toda lattice vertex operator (see [2]), given by

$$\begin{aligned}\mathbf{X}(t, y, z) &= -\chi^*(z)X(-t, z)X(t, y)\chi(y) \\ &= \frac{z}{y-z} e^{\sum_1^\infty t_i(y^i-z^i)} e^{-\sum_1^\infty (y^{-i}-z^{-i})\frac{1}{i}\frac{\partial}{\partial t_i}} \left(\frac{y^n}{z^n} \right)_{n \in \mathbf{Z}};\end{aligned}$$

in particular, if τ is a 1-Toda vector, then $a\tau + b\mathbf{X}(t, y, z)\tau$ is a 1-Toda vector as well. Using the dictionary, this leads to q -vertex operators

$$\mathbf{X}_q(x, t; y, z) = e_q^{xy} (e_q^{xz})^{-1} e^{\sum t_i(y^i-z^i)} e^{-\sum (y^{-i}-z^{-i})\frac{1}{i}\frac{\partial}{\partial t_i}} \quad \text{for } q\text{-KP},$$

and, for any N th root ω of 1,

$$\mathbf{X}_q(x, t; z) = e_q^{x\omega z} (e_q^{xz})^{-1} e^{\sum t_i z^i (\omega^i - 1)} e^{-\sum z^{-i} (\omega^{-i} - 1) \frac{1}{i} \frac{\partial}{\partial t_i}} \quad \text{for } q\text{-KdV},$$

having the typical vertex operator properties.

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