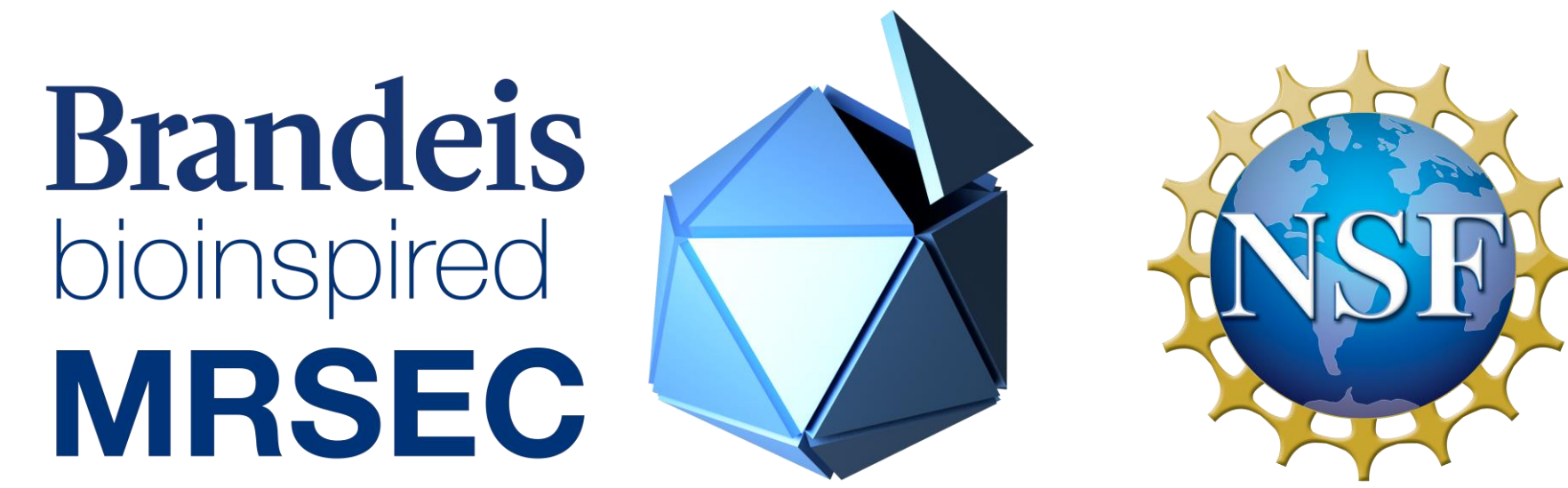


Instabilities of 3D dry active nematics: particle-based simulation and perturbation analysis

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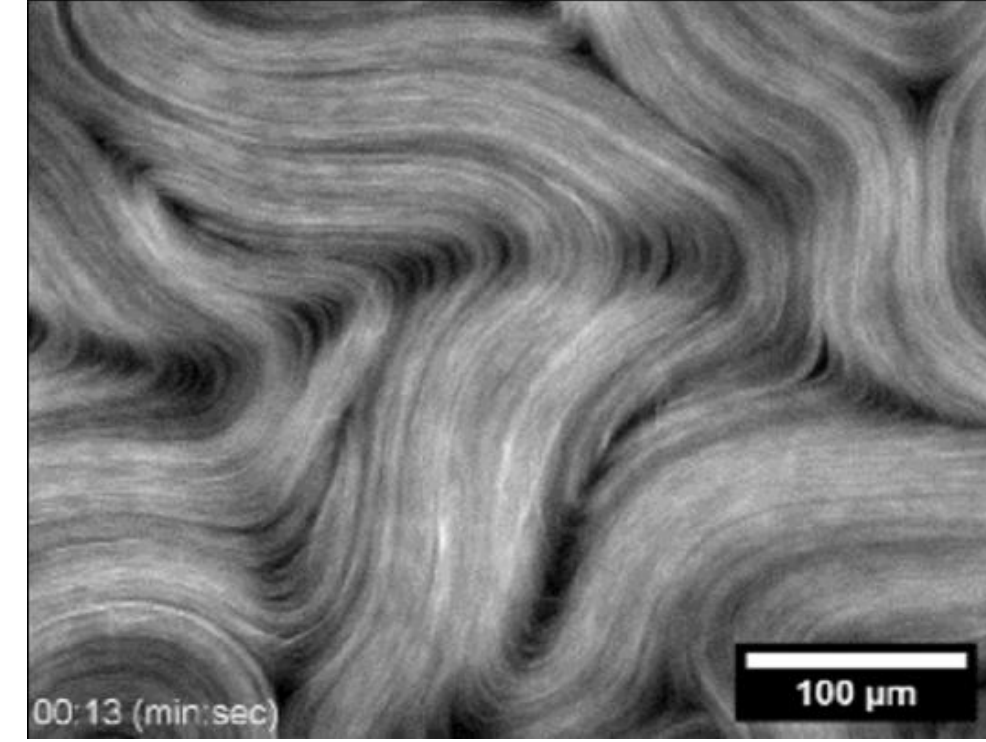
Brandeis University, Martin A. Fisher School of Physics



Background

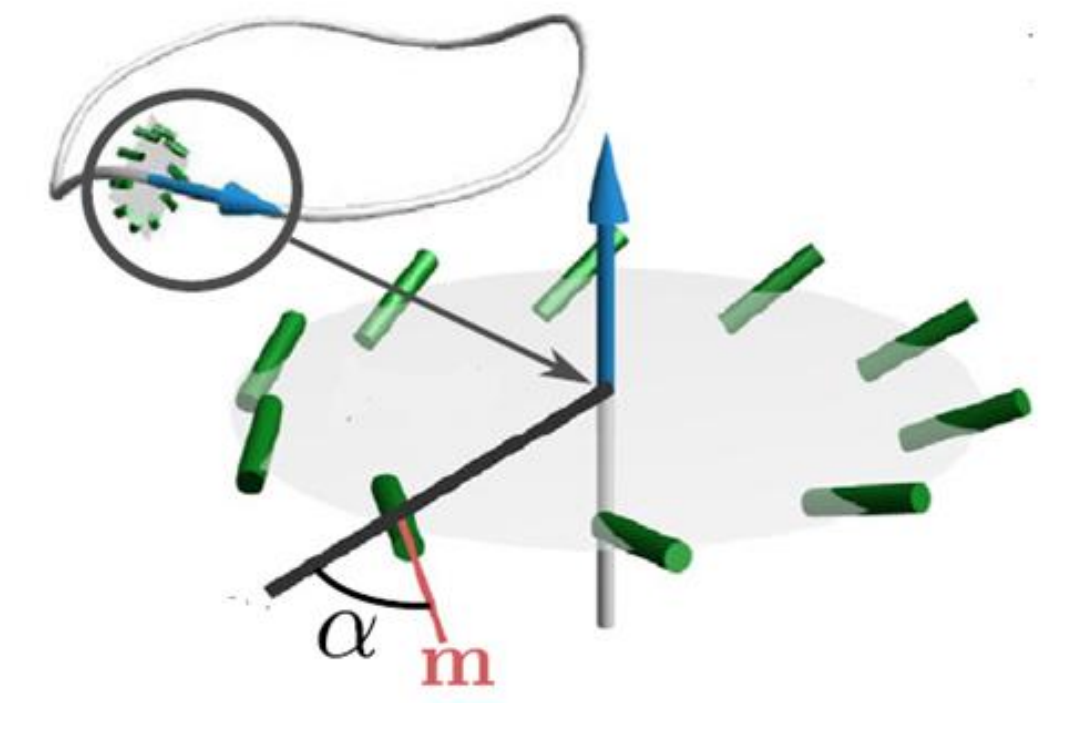
With self-propelled velocity, nematics may deviate from aligned ordered state to form pairs of topological defects. In 3D, such defects will be line-like and the simplest model is disclination loop. To analyze the location and density of defects can strongly help us understand the configuration and dynamics of active nematics.

2D active nematics



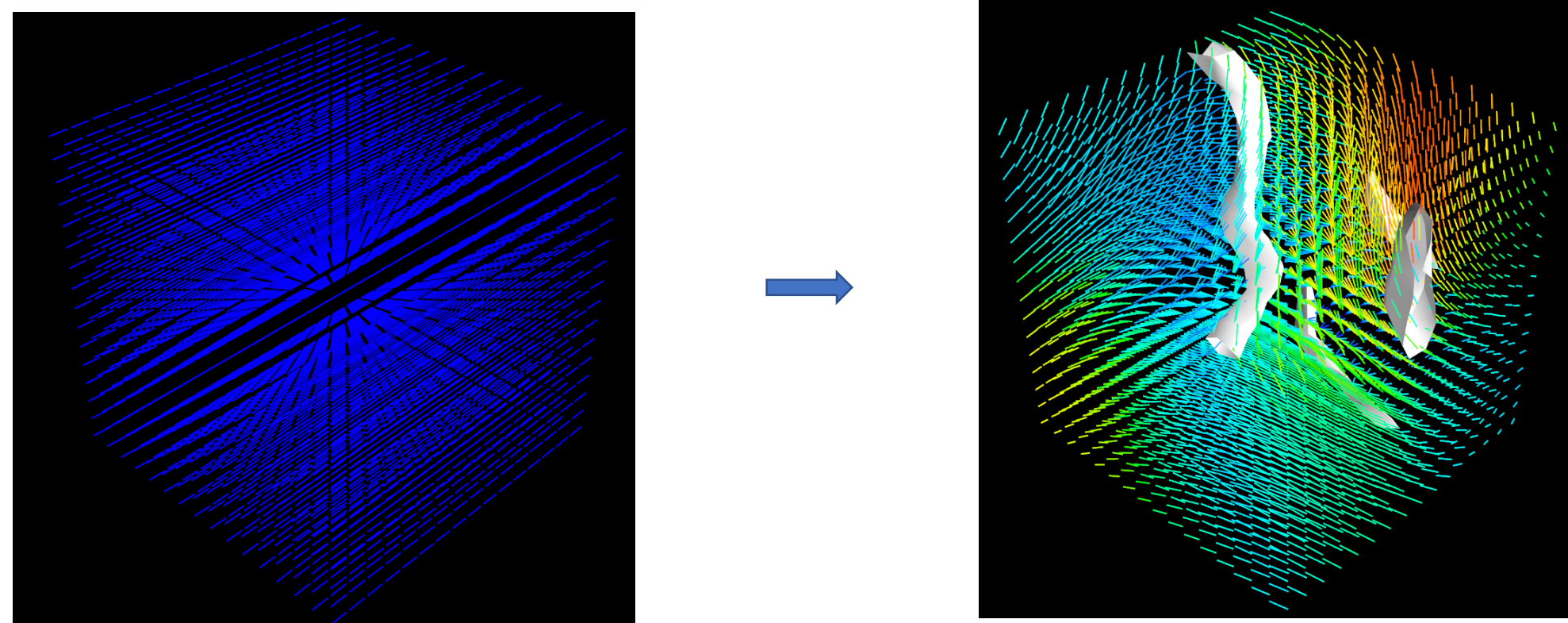
DeCamp, et al, Nature, 2015

3D disclination loop



J. Binysh, et al, Phys. Rev. Lett., 2020

Instability



Colored dash: director and its orientation

White region: disordered lines

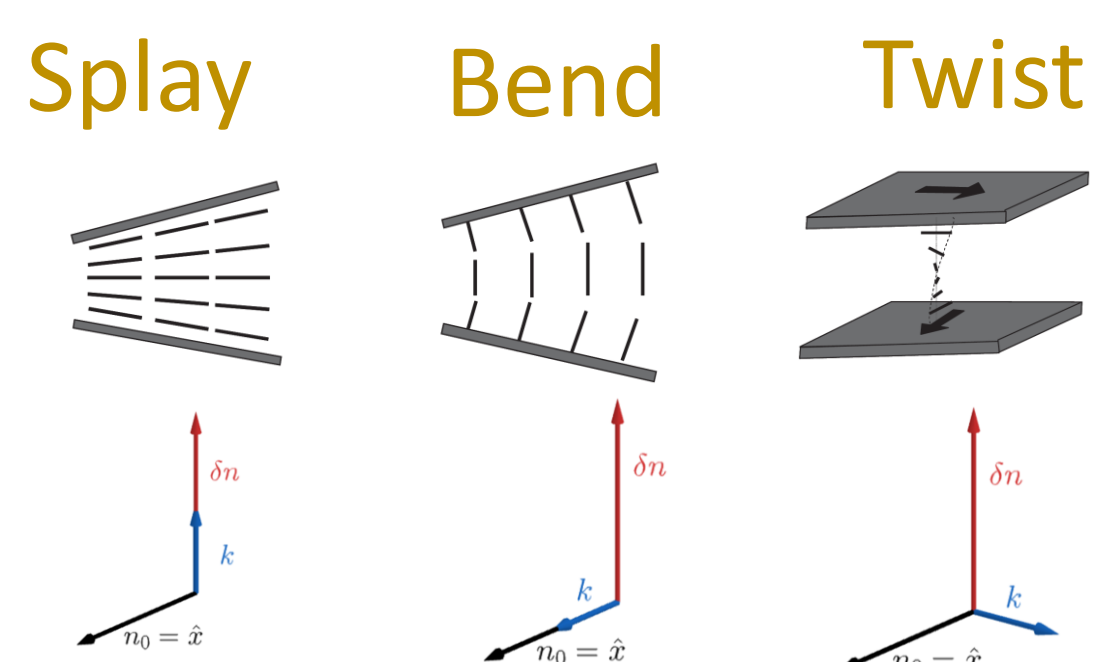
Introduction

Unlike the previous research, in this project we focused on the **DRY** system, namely there is no hydrodynamics interaction or no solvent. This can help us understand how only active force and filaments' alignment influence the emergent dynamics.

We aimed to investigate the **instability** of 3D active nematics: how the filaments reorient from initial aligned state and form disclination loops.

Question 1: What deformation appears first? By perturbation analysis on continuum description

There are 3 different kinds of deformation modes in nematics. With δn (linear perturbation) and Fourier transform, we can depict each mode by orientation of wave vector k .



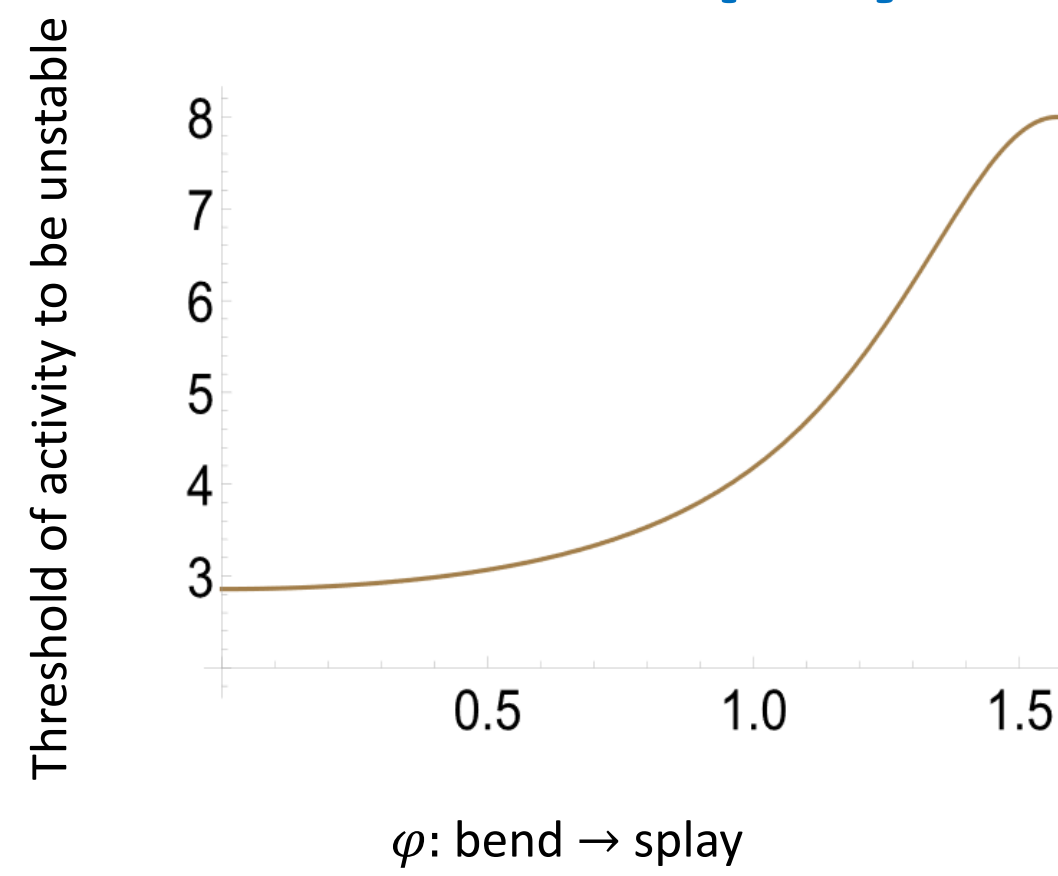
Also by such perturbation analysis, we investigate the equation in the **right**, to find which orientation of k is the most unstable.

$$\partial_t \mathbb{Q} + \mathbf{u} \cdot (\nabla \mathbb{Q}) = -\frac{1}{\gamma} \frac{\delta \mathcal{F}}{\delta \mathbb{Q}} + (\mathbb{Q} \Omega - \Omega \mathbb{Q}) + \lambda \left[\mathbb{Q} \mathbf{E} + \mathbf{E} \mathbb{Q} - \frac{2}{3} \text{Tr}(\mathbf{E} \mathbb{Q}) \mathbb{I} \right] + \lambda \rho \left[\mathbf{E} - \frac{1}{3} \text{Tr}(\mathbf{E}) \mathbb{I} \right]$$

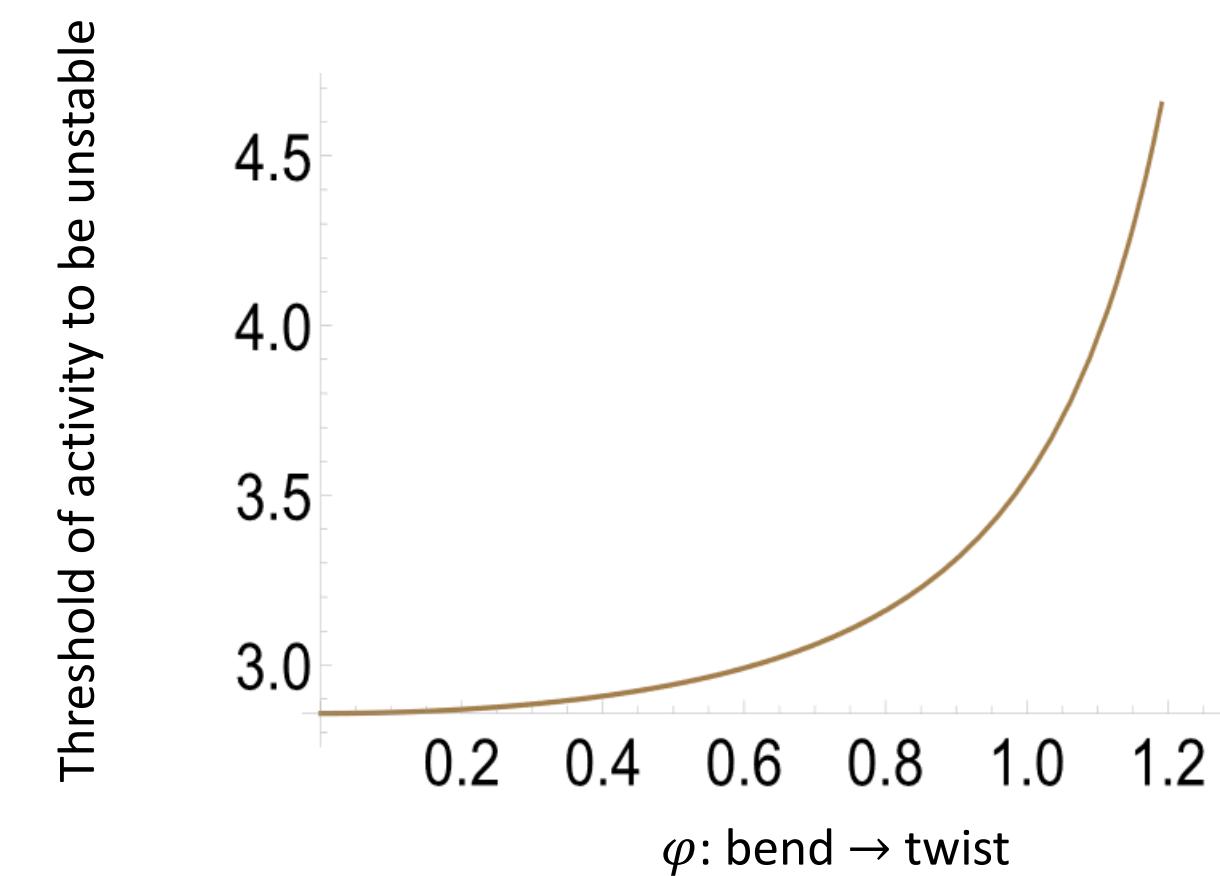
$-\zeta \mathbf{u} = \alpha \nabla \cdot \mathbb{Q}$ **DRY system:** velocity only comes from activity

Considered the system of two coupled deformations for simplicity:

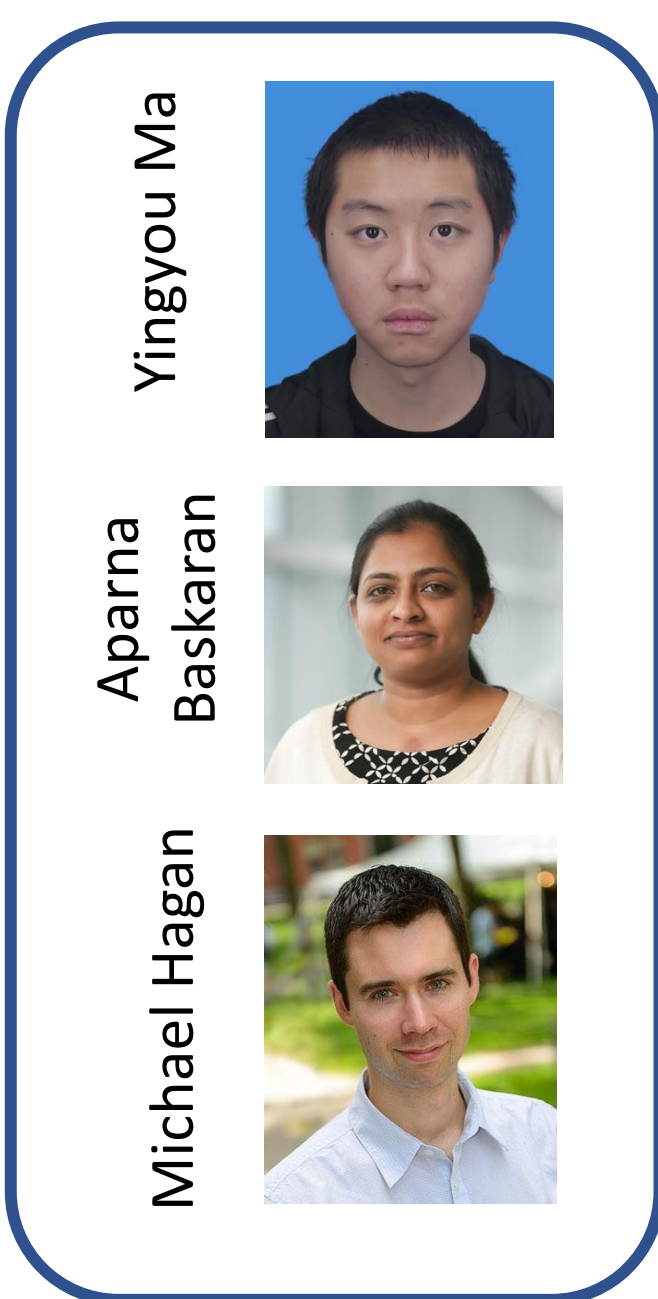
Bend-Splay



Bend-Twist



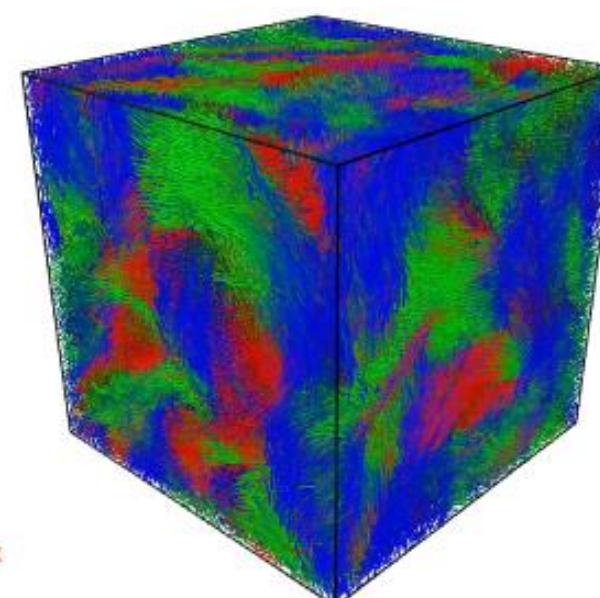
Conclusion: Bend will firstly go unstable



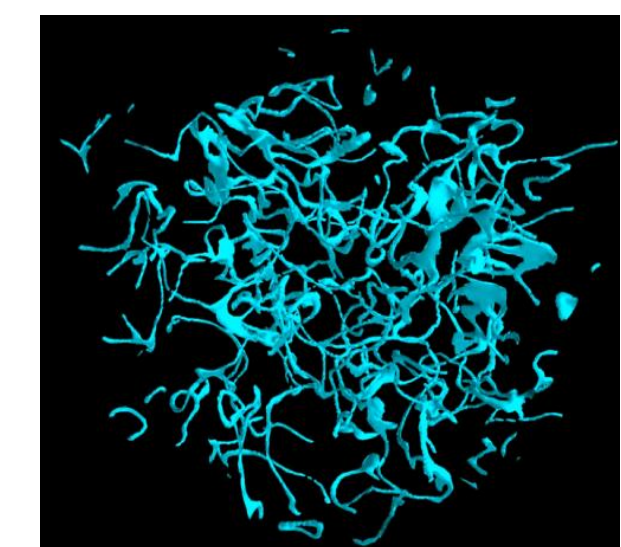
Question 2: How disclination loops emerge? By particle-based simulation (semi-flexible filaments)

$$m_i \ddot{\mathbf{r}}_i = \mathbf{f}_i^a - \gamma \dot{\mathbf{r}}_i - \nabla_{\mathbf{r}_i} U(\mathbf{r}_{ij}) + \mathbf{R}_i(t)$$

Activity Friction Potential Thermal Noise

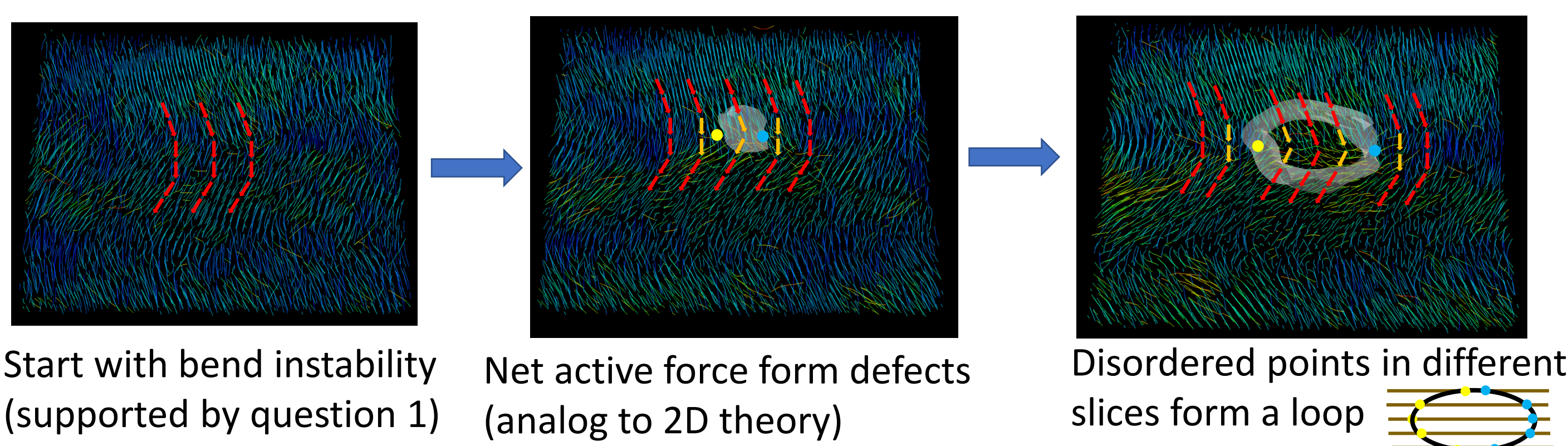


DRY system:
Only particles of nematics



The network of disordered regions

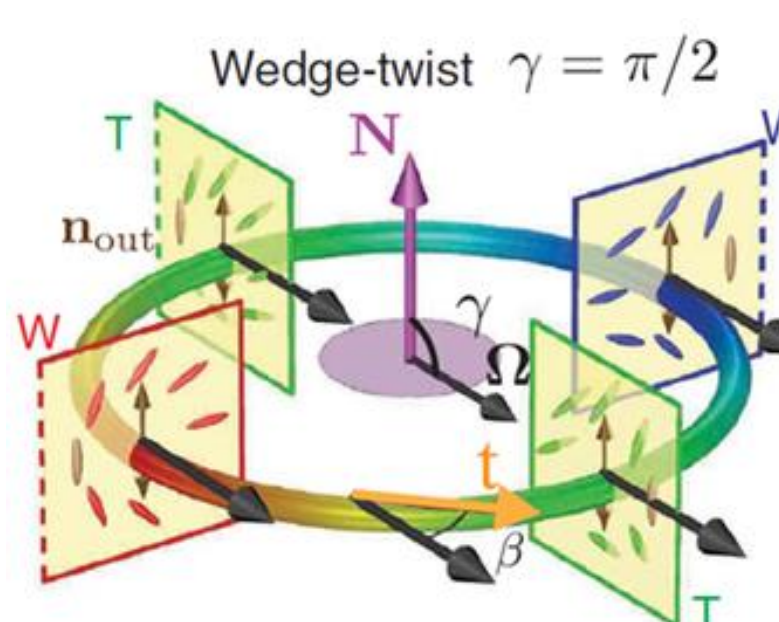
Based on 2D theory of forming defects, we derive a hypothesis of the 3D disclination loops' appearing (white: disordered):



Start with bend instability (supported by question 1)

Net active force form defects (analog to 2D theory)

Disordered points in different slices form a loop



This hypothesis will always form wedge-twist loops. In the shown trial, every loop is wedge-twist.

G. Duclos, et al, Science, 2020

Conclusion: 2D-analog hypothesis

Summary

We investigated the initial instability period of active dry nematic. The perturbation analysis shows bend is always firstly unstable. A mechanism of forming a disclination loop was proposed.

Future work

Check if there is another way of forming a loop. Investigate the evolving and interaction of loops. Check the results by continuum simulation.