

Force as a variable to study unfolding reactions

- Many biological reactions proceed by large conformational changes that may be force dependent
- By "stretching" a molecule ~~is-a~~ along a particular direction, we can study force induced conformational changes along a well defined mechanical reaction co-ordinate

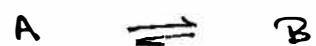
↓
These can be extension, position of a motor along its track etc

- These experiments often yield information about the free energy landscape of a reaction. Some of the obtainable information are 1) ΔG 2) distance to the transition state & 3) number of intermediates

Illustrative example: unfolding of a RNA hairpin

Liphardt et al. Science 2001 Apr 27; 292(5517)
733-737

Simple two-state system

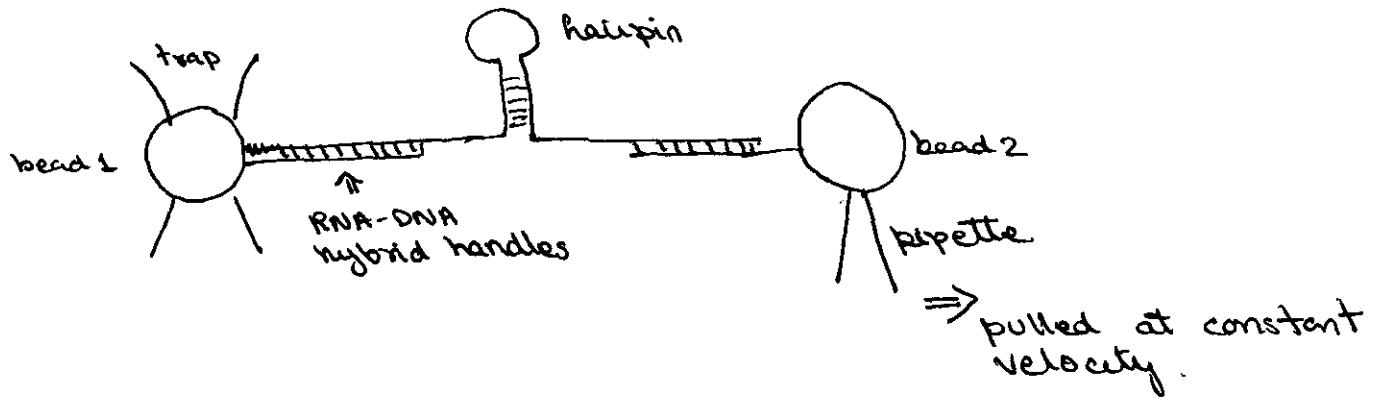


mechanical co-ord: end-to-end distance

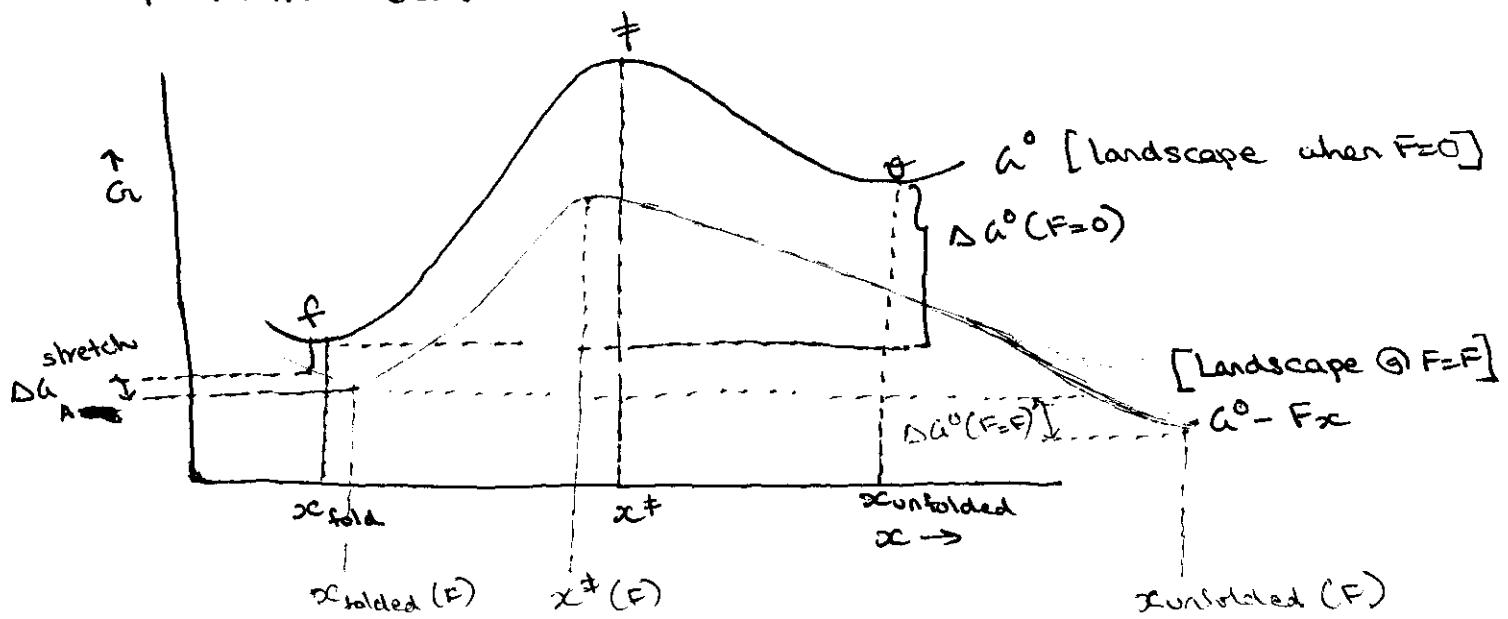
(measured as dist
between the two
beads)

(see next pg for
end set-up)

Exp. design



How does force affect the landscape of an unfolding reaction? \Rightarrow A theoretical discussion before looking at experimental data



- Effect of force on ΔG :

$$\Delta G_{A \rightarrow B}(F) = \Delta G_{A \rightarrow B}(F=0) - F(x_B - x_A) + \Delta G_{A \rightarrow B}(\text{stretching})$$

$A \equiv$ folded
 $B \equiv$ unfolded

where, $\Delta G_{A \rightarrow B}(F=0) = \Delta G_{A \rightarrow B}^0 + k_B T \ln \frac{[B]}{[A]}$

$\Delta G_{A \rightarrow B}^{\text{stretch}}$ arises from the shifting of the minima of the folded states to the right if is given by :-

$$\Delta G_{A \rightarrow B}^{\text{stretch}}(F) = \left[\int_{x_B(F)}^{x_A(F)} F dx - \int_{x_B(0)}^{x_A(0)} F dx \right]$$

Summary points :-

- Force "tilts" the free energy landscape of a reaction by $F \cdot x$

$$G(F) = G(0) - F \cdot x$$

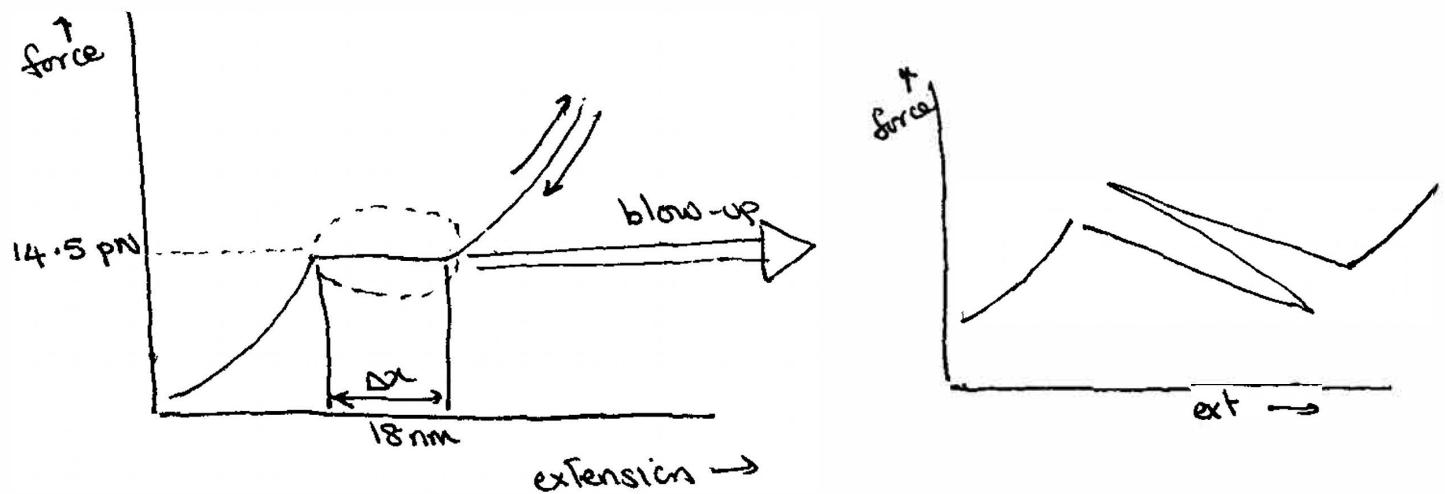
- This results in the unfolded state being stabilized greater than the folded state
- A more extended molecule is stabilized with force such that $\Delta x_{A-B}(F) > \Delta x_{A-B}(F=0)$
- Rate of the unfolding reaction increases as the transition state is stabilized more relative to the folded state
- At equilibrium $\Delta G_{A \rightarrow B}^{\text{(F)}} = 0$

$$\Delta G_{A \rightarrow B}^{\text{(F)}} = \Delta G^\circ - F \Delta x + k_B T \ln \frac{[x_B]}{[A]} + \Delta G_{A \rightarrow B}^{\text{stretch}}(F) = 0$$

$$\Delta G^\circ = F \Delta x - k_B T \ln \frac{[B]}{[A]} + \Delta G_{A \rightarrow B}^{\text{stretch}}(F)$$

Experimental results from RNA hairpin unfolding experiments :-

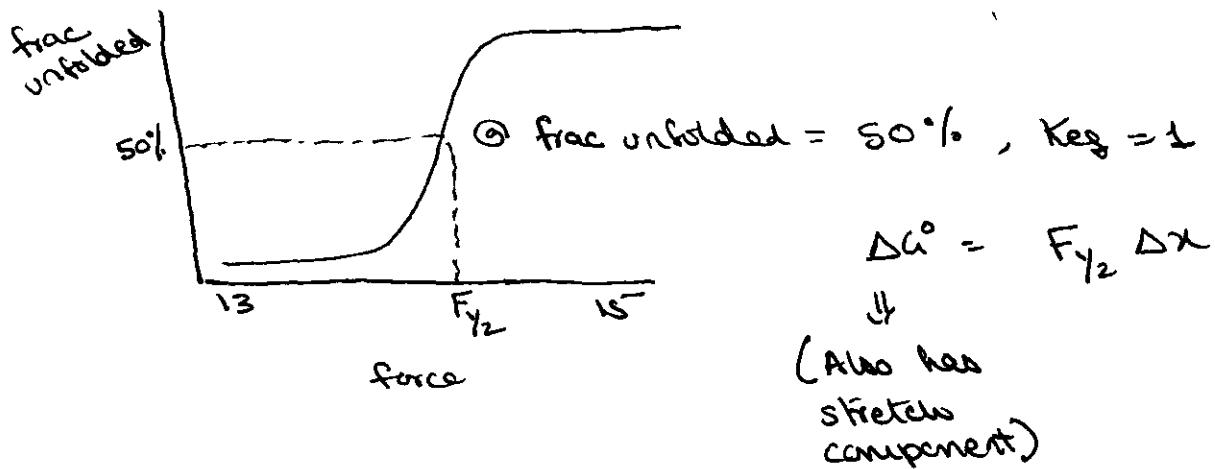
- Experiment in which the micropipette attached bead is pulled at a constant velocity thus increasing the force on the molecule at a constant rate



- stretching & relaxing curves superimpose indicating that the reaction is reversible & happening at equilibrium
- Because we are seeing a process at equilibrium, "hopping" between folded & unfolded states is observed at the unfolding transition plateau.
- Average area under the unfolding plateau

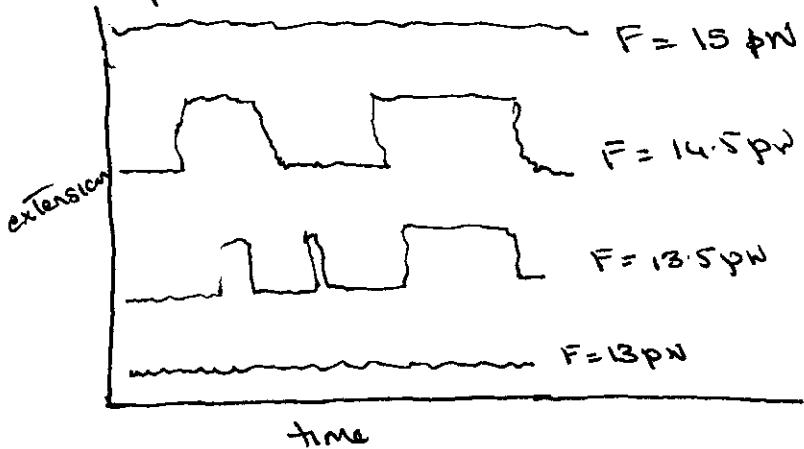
$$W = \int_{x_0}^{x_0 + \Delta x} F \cdot dx = \Delta G^\circ$$

- Another way to replot the data



$\Delta G^\circ \approx 150 \text{ kJ/mol}$, compared well with predicted value

- Experiment at constant force



$$K_{eq} = \frac{k_{fold, B \rightarrow A}}{k_{unfold, A \rightarrow B}}$$

$$k_{unfold, A \rightarrow B}(F) = k_0 e^{-F\Delta x^+ / k_B T}; \quad \Delta x^+ \rightarrow \text{dist to transition state}$$

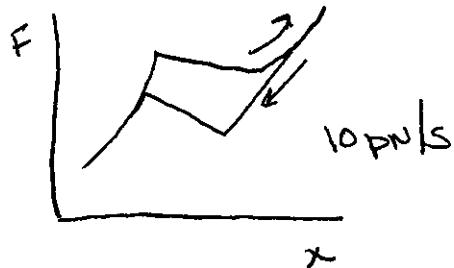
$k_0 = \text{rate constant at } F=0$

- Reactions not occurring at equilibrium.

For the hairpin with a bulge in the structure



x
reversible, equilibrium
process. $\Delta H = W_{\text{rev}}$



x
 r^* not at equilibrium,
not reversible

$$W_{\text{irrev}} > \Delta H$$

Hysteresis indicates that the reaction is irreversible \leftarrow not occurring at equilibrium. This results when the loading rate ($r = dF/dt$) is greater than the slowest relaxation rate of the system

Unfolding force depends on loading rate

\rightarrow higher loading rate \equiv higher unfolding force

as the time spent Θ each force is short

Probability of unfolding as a function of force + loading rate is given by:-

$$P_u(F, r) = 1 - e^{-\frac{k_B k_B T}{r \Delta x^2} (e^{F \Delta x^2 / k_B T} - 1)}$$

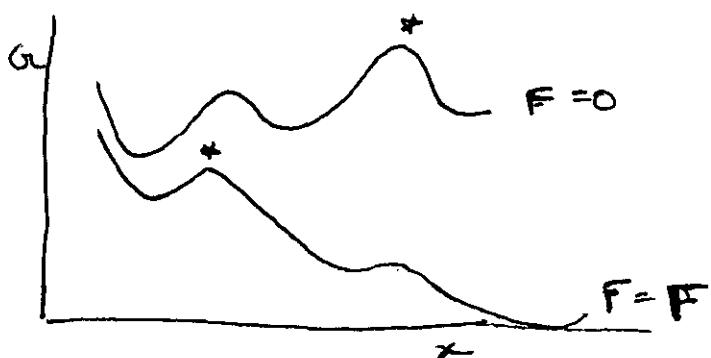
where, Δx^\ddagger = distance to the transition state
 k_0 = unfolding rate constant at zero force

The most probable unfolding force is given by:-

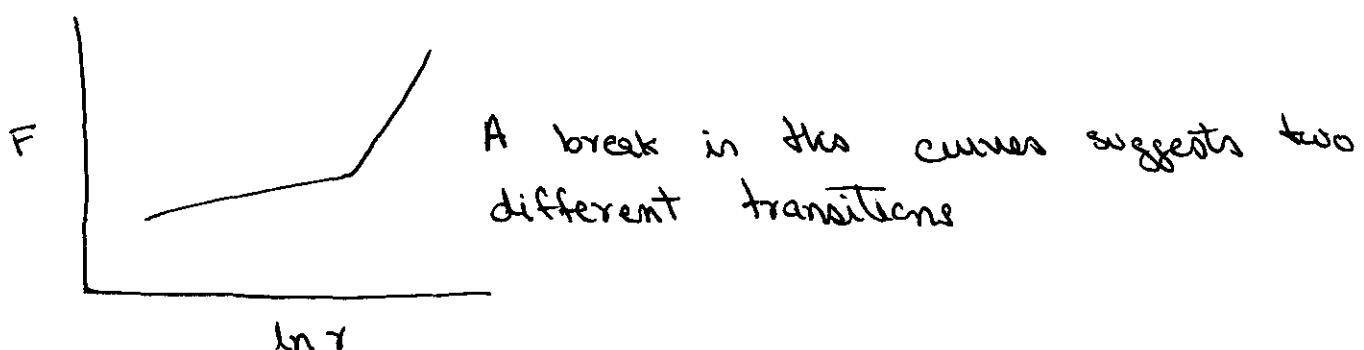
$$F^* = \frac{RBT}{\Delta x^\ddagger} \ln \left(\frac{\gamma \Delta x^\ddagger}{k_0 e^{-F^*}} \right)$$



For a complex landscape with multiple energy barriers



* significant energy barrier can change when reaction is occurring at different forces



This illustrates why "k₀" is not a very meaningful quantity as the intercept of the two curves give different values neither of which may be the true "k₀" at F=0

Discussion points:-

- Mechanical versus chemical unfolding results
 - Thermodynamic quantities like ΔG should be comparable
 - Kinetic τ quantities are path dependent & should be compared with a lot of caution.
- Mechanical ~~verses~~ versus thermodynamic stabilities
 - \downarrow F^*
most probable unbinding force
 - \uparrow ΔG
free energy
 - These are NOT the same quantity
 - Mechanical stability depends on both the rate of pulling & the direction of pulling
 - Mechanical stability is the response of a molecule to force & depends on the distance to the transition state. This information is not contained in ΔG .
- Biological relevance of mechanical stability.
 - Unfolding processes & rupture processes in biology are probably better described by mechanical rather than thermodynamic stabilities. However, it is important to consider direction/pulling rates in-vivo to make the measurement more biologically relevant.